

**FEA Validation of NiTiPdPt and NiTiPdCu Shape Memory  
Alloys at High Temperatures**



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August, 2016**

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A thesis submitted in partial fulfilment of the requirements for the degree of  
MS Design and Manufacturing Engineering

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Design and Manufacturing Engineering  
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August, 2016

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I certify that this research work titled “*FEA Validation of NiTiPdPt and NiTiPdCu Shape Memory Alloys at High Temperatures*” is my own work. The work has not been presented elsewhere for assessment. The material that has been used from other sources it has been properly acknowledged / referred.

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## **Dedication**

*I dedicate my efforts to my beloved parents who made me what I am today, my teachers for guiding me all the time consistently, my siblings for backing me up in all thick and thin and my students for inspiring me.*

## **Abstract**

Shape Memory Alloy (SMA) also known as smart metal, memory metal, memory alloy, muscle wire, smart alloy is an alloy that "remembers" its original, cold-forged shape: returning to the pre-deformed shape when heated. This material is a lightweight, solid-state alternative to conventional actuators such as hydraulic, pneumatic, and motor based systems.

The Cu-based, Fe-based and NiTi-based SMAs are considered as engineering materials. The yield strength of SMA is lower than that of conventional steel, but higher yield strength than plastic or aluminium. The yield stress for NiTi can reach 500 MPa.

Some devices are made to test the properties of SMA but the research is still in progress to authenticate the experimental results with the true behaviour of SMA. To cope up with this problem, a numerical simulation based solution is needed to be devised to check the thermomechanical properties of newly constituted SMA and verify its load and temperature bearing capacity as well as deformation evaluation.

In this research, one experimentally tested and verified SMA NiTiPdPt is simulated Finite Element Analysis package and the results are compared with the newly constituted SMA NiTiPdCu. The results found for both materials verify the authenticity and validation of FEA method for evaluation of thermomechanical properties of SMA.



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## Nomenclature

<b>Symbol</b>	<b>Description</b>	<b>Unit</b>
$A$	Cross-sectional area	$\text{m}^2$
$\alpha$	End-condition constant	
$C$	Spring index	
$C'_1$	Elastic constant	
$C'_2$	Elastic constant	
$d$	Diameter of spring wire	$\text{m}$
$D$	Mean diameter of spring coil	$\text{m}$
$E$	Modulus of elasticity	$\text{GPa}$
$F$	Axial force	$\text{N}$
$G$	Modulus of rigidity	$\text{GPa}$
$J$	Mass moment of inertia	$\text{m}^4$
$k$	Spring rate or scale	$\text{N m}^{-1}$
$K_W$	Wahl factor	
$K_s$	Shear-correction factor	
$K_B$	Bergsträsser factor	
$L_0$	Free length of spring	$\text{m}$
$\lambda_{\text{eff}}$	Effective slenderness ratio	
$N$	Number of coils	



$N_a$	Number of active coils	
$r$	Mean radius of spring coil	m
$S_{ut}$	Ultimate tensile strength	MPa
$S_{sy}$	Torsional yield strength	MPa
$S_y$	Yield strength	MPa
$T$	Torsional moment	J
$\tau$	Shear stress	Pa
$\tau_{\max}$	Maximum shear stress	Pa
$U$	Total strain energy	J
$y$	Deflection	m
$y_{cr}$	Critical deflection	m



# CHAPTER 1

## INTRODUCTION

### 1.1 Significance

There is a number of metallic alloys that have been exhibit the shape memory effect. Wayman gives a comprehensive definition of alloy (SMA) which involves the majority of all the thermal shape memory alloys. He states

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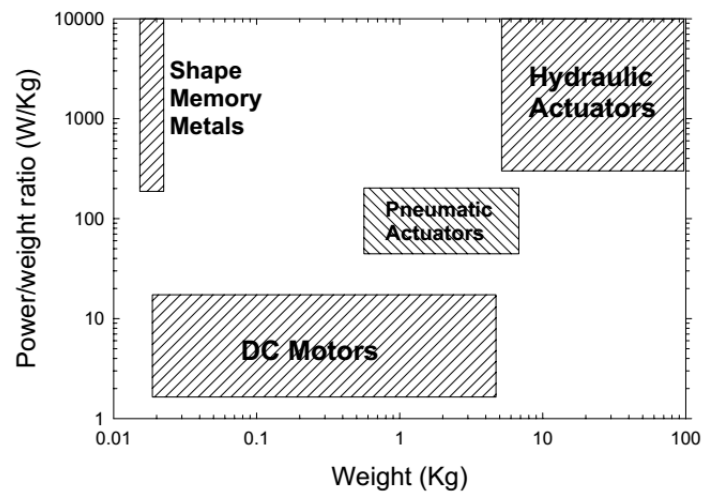
“A thermal shape memory alloy as an article that when deformed at a lower temperature will regain its *original* shape when heated to a higher temperature [1].”

---

Shape memory alloys are utilized as a part of numerous engineering applications [2]. The most common commercial system is NiTi based SMAs. Applications up to this point for NiTi combinations include electrical switches, micro-actuators, temperature sensitive valves, appliance controllers, cell phones antennas, brassiere underwires, eyeglass frames and countless dental and medical devices [3, 4]. In addition, the first vast scale commercial applications for shape memory alloys were made utilizing NiTi and NiNbTi alloys with sub-rooms changed temperatures, for use as coupling for tubes, funnels, tubes, and electrical interconnects [5]. These applications make utilization of NiTi alloys close room temperature. The fundamental reason that commercial applications have been constrained to close room temperature is that commercial NiTi SMAs have a highest changed temperature of about 100 °C.

Likewise, there are many control and activation-sort applications for materials displaying the shape memory impact at maximum temperatures. Maximum temperature alloys (HTSMA) could be utilized as a part of an aeronautic, power generation [6], chemical processing industries and automotive. While particular applications have been distinguished in light of some type of a HTSMA, no appropriate material have been created. As is normal in the material the improvement of utilizations for cutting edges is somewhat in front of material advancement

itself, and such is the situation for the advancement of high- temperature shape memory compounds.



**Figure 1.1** Power-to-weight ratio versus weight diagram for common actuator types currently used in aeronautics [6]

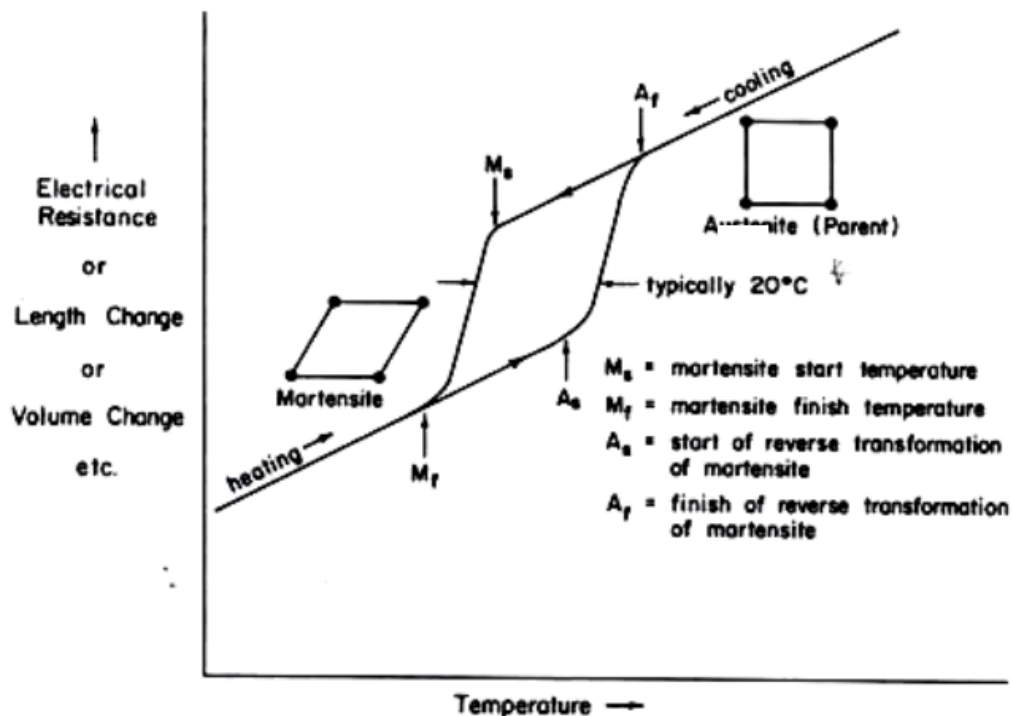
Integration of SMA actuators into aeronautic turbomachinery would bring about few natural advantages. Aeronautics plainly stresses weight reduction in all stages of engineering. Decreases in the net weight brings about sizable expansions in fuel effectiveness. Also, SMA actuators diminish the quantity of subsystems when contrasted with standard with typical pneumatic or more common hydraulic and motor-driven actuators, giving additional decreases in weight and cost. Figure 1-1 demonstrates the common weight to power proportions of the more typical commercial actuators now utilized by the aviation industry [7]. Limiting weight and expending power brings about an execution file in which SMA actuators are plainly predominant.

The design and development of effectively controlled SMA gadgets requires top to bottom portrayal of the mechanical and shape memory particular properties. Past reviews have been represented compositional impacts of change temperature and, in a few cases, stack free recuperation, yet there is an entire absence of information required for the application of shape memory alloys especially in actuator-related applications. This study endeavors to describe these properties and associate them to material composition and microstructure, which thusly can be utilized to distinguish conceivable ranges for further alloy and process development.

## 1.2 Background

### 1.2.1 General Shape and Memory Alloy Behavior

SMA's are categorized by a set of temperatures at which a crystallographic auxiliary change starts and stops. The high-temperature austenite or parent stage is a high symmetry stage generally requested while the lower temperature martensite stage is a lower symmetry structure which shapes from the high symmetry parent stage diffusion-less shear change. The different temperature at which this change starts and closures on warming and cooling the characterized  $A_s$ ,  $M_s$ ,  $A_f$  and  $M_f$ .



**Figure 1.2** Power-to-weight ratio versus weight diagram for common actuator types currently used in aeronautics [8].

The austenite begin temperature,  $A_s$ , is the temperature at which the change of the martensite to austenite stage starts on warming.  $A_f$  is the temperature at which the change is finished and the material is 100% austenite. The martensite start,  $M_s$  and martensite end,  $M_f$ , temperatures are the temperatures at which the change happens on cooling [9].

Figure 1.2 is a traditional schematic introduction by Wayman which demonstrate a material property subordinate changes a shape memory alloy is burned through a thermal hysteresis [8].

An Irregularity in the materials properties emerges at the beginning of the change in heating or cooling, which is typical for all SMA materials. All through the change the material properties, this property might be for instance particular volume, electrical resistivity, modulus, or other fundamentally subordinate property. The most widely recognized test strategies for the assurance of the change temperatures are thermal techniques (DSC and DTA) dilatometric strategies, and resistive techniques. The last two outcome in comparable hysteresis plots as exemplified in Figure 1.2. While this figure 4 is glorified it is illustrative of a large number of the significant qualities of thermal shape memory alloys.

### 1.2.2 SMA Structural Characteristics

Shape memory alloys display both a thermodynamically and crystallographically reversible change. A crystallographically change is probably at the point when the interface between the martensite and austenite is basically cognizant and the parent austenite stage is an ordered compound. The high-temperature austenite stage is a higher-symmetry structure, normally ordered cubic ( $B2$  structure) as in case of NiTi, also, NiTi changed compounds. The austenite stage changes without considerable long range dispersion into a lower symmetry martensite structure at some decrease temperature. Basic cells schematically delineating the structures basic in NiTi higher and lower the change temperatures are appeared in Figure 1.3 [10].

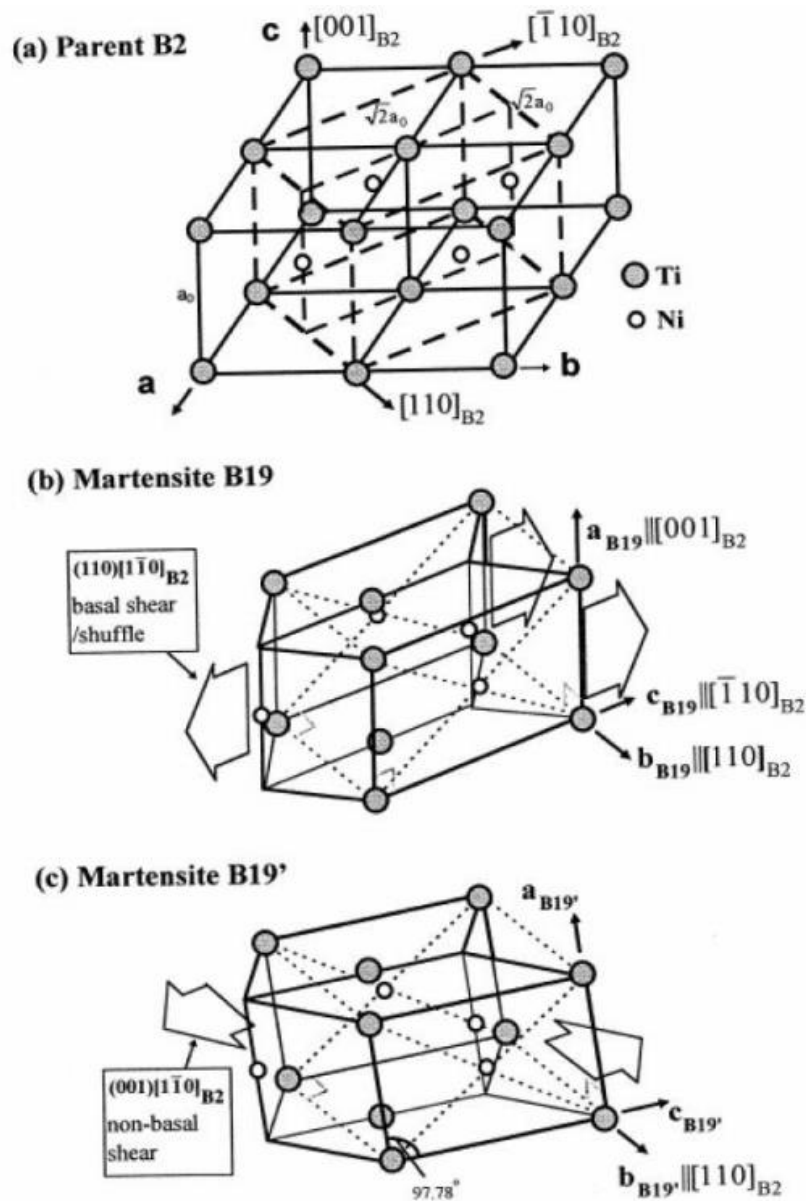
The cubic  $B2$  parent stage in NiTi based SMAs changes to various distinctive martensitic structures. The last structure lies on upon alloying additions, polluting influences, and preparing history [7, 9]. The following change responses have been distinguished.



Every response is crystallography reversible except for the  $B2 \rightarrow R$  change. The  $R$  stage is accomplished by stretching of the  $B2$  structure resulting in a rhombohedral structure [10].

The  $B19$  stage is an orthorhombic structure which is shaped from the  $B2$  parent crystal in a few stages, which comprise of prolongation about the  $a$ ,  $b$  and  $c$  axes and shearing of the basal plane in the  $c$ -direction which is common to the  $b$  direction. The  $B19$  non-basal plane, which

is typical to the  $a$ -bearing in the  $c$ -direction [10]. The  $B19'$  stage is the essential shear structure which show up in parallel NiTi leading to the shape memory impact.



**Figure 1.3** Structure of the parent phase (austenite) and shear phases (B19 and B19' martensite) [10].

The cubic parent stage in NiTi SMA is steady at high temperature (over the AF). Numerous lattice variations of martensite can shape from each parent austenite grain. Each variant will take after a superbly reversible way back to the parent austenite stage due to the ordered way of the compound. On the off Chance that this did not happen, the material would experience a diffusional change in the synthetic ordering of the first parent stage's lattice [11].

In the Pt and Pd adjusted NiTi SMAs the  $B19$  stage is of essential significance since alloying increases more noteworthy than around 10% outcomes in the essential martensitic structure changing from the monoclinic  $B19'$  to the orthorhombic  $B19$  [12]. The  $B19$  and  $B19'$  are shear structures of the  $B2$  and in this way, display a lower symmetry, with the  $B19'$  being the shear structure of the  $B2$  and in this way, display a lower symmetry, with the  $B19'$  having the least symmetry [13]. The symmetry of the structure is of significance as it is the fundamental consideration of the assurance of the quantity of identical martensitic variations which may frame from a parent  $B2$  cubic structure. The  $B19$  has 12 proportional variations that may frame from a parent crystal. This means 12 distinctive approaches to shear the  $B2$  structure in the development of the  $B19$ . Every identical  $B19$  structure may then bearing to shape a  $B19'$  martensite which brings about 24 comparable variants [14]. By and large when breaking down shear structures the quantity of proportionate variations increments as the symmetry of the shear structures the quantity of proportionate variations increments as the symmetry of the shear structure diminishes. This relationship will be analyzed in further detail as it identifies with the disfigurement behavior and mechanisms in the shape memory alloy covered in this study.

The primary theory of martensite development, which applies to both thermoplastic and thermoelastic changes, relays a high symmetry parent structure to a variation of the shear structure. Despite the fact that this theory is appropriate to many shear stage changes for graphic purposes that will concentrate in the structures appeared in Figure 1.3. Three key disfigurement steps are essential to identify with structures by an absolutely shear transformation [15].

1. Primarily there is a Bain distortion which achieves its name generally from the distortion seen in the thermoplastic change which shapes the metastable tetragonal martensite. The Bain distortion is basically the lengthening of the parent stage which is appeared in Figure 1.3 as the extensions of the  $a$ ,  $b$  and  $c$  axes.
2. Secondly a shear distortion must happen so as to safeguard the grid symmetry which joined with the Bain distortion frames the undistorted plane or propensity plane.
3. At long last there is a pivot which brings the undistorted plane into a similar orientation in both the parent and shear stage.

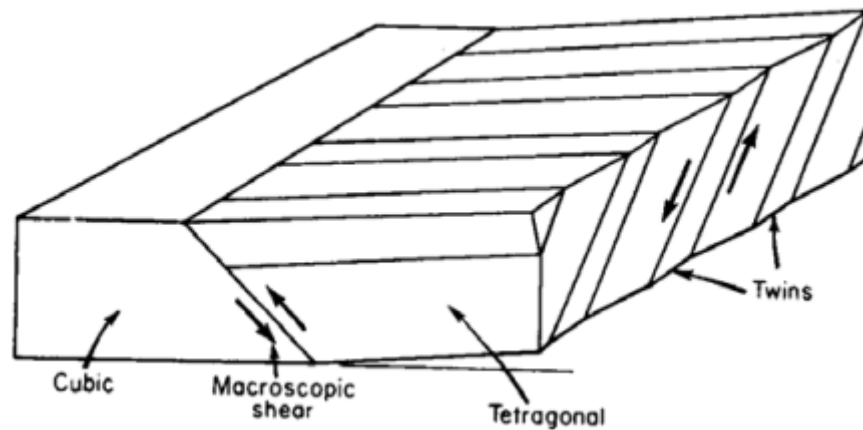
The accommodation of random shearing of the lattice (noted as step 2) is a key component characterizing whether a shear change is thermoplastic or thermoelastic. Thermoplastic



changes, for example, those moral in steels are non-reversible and the greater part of the change shear related with step 2 is accommodated plastically. That is, there is the development of non-reversible imperfections, for example separation movements and generation. Thermoelastic changes then again accommodate the greater part of the change shear flexibly in mix with recoverable mechanisms.

There are three fundamental ways the change shear might be suited, two of which are adjustable. The dynamic shear component depends to a limited extent on the mechanical properties of the austenite and martensite including the magnitude of the change shear. In both thermoelastic and thermoplastic changes a lot of shear is accommodated flexibly which is dependent to the yield quality of the martensite and austenite nearby the interface [16]. If the change strain brings about an interface stretch state which surpasses its nearby yield quality change strains are suited irreversibly by plastic distortion or slip. At long last, if the shear stress at the interface required to start disfigurement twins is adequately low and higher than the yield quality the change strain is suited by twin formation [17].

Twin accommodated strain, which is the foundation of shape memory alloy, forms during the change along the twin planes in the shear (martensite) stage. It is critical to note that the twinning plane is generally a low index plane that is not parallel to the habit plane. Figure 1.3 is a schematic of the way of a thermoelastic change interface amongst parent and shear stages in which the change strain is accommodated by twin arrangement. The microscopic visible shear plane, which isolates the cubic and shear structures is subject to the auxiliary connections between the martensite and austenite which incorporate Bain shear and pivot [18]. The twinning plane however depends on the symmetry of the shear structure as the deformation twin must just reorient the structure by an associated development of molecules consistently conveyed over the volume isolated by the twinning plane. The change shear along the plainly visible shear plan is in part accommodated by the shear related with twinning along the twin planes [19]. In thermoelastic changes the net changes shear in the martensite is suited to some degree by the versatile deformation of the twinned and untwined locales and incompletely by the formation of the deformation twins. If the deformation twins did not arise the pressure because of the changes strain would surpass the yield quality of the material and in this manner, be accommodated plastically. Twins that shape through thermoelastic changes, with a specific end goal to accommodate the shear along the plainly visible shear plane, have a twin linked crystallographic relationship to each other [20].



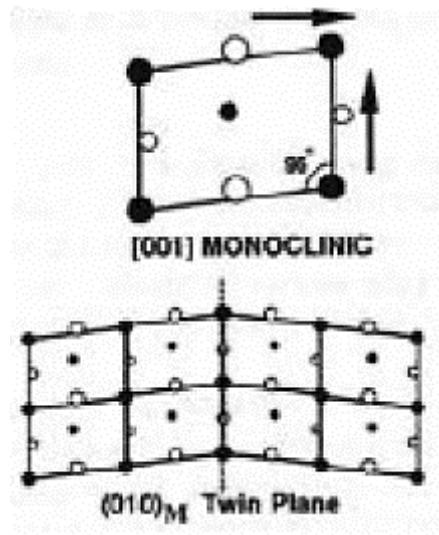
**Figure 1.4** Thermoelastic transformation and twin accommodated transformation strain [15].

As exemplified in Figure 1.4 the feeling of the shear related with a twin must substitute between twinned areas. This rule for the settlement of the change strain brings about the arrangement of coupled sets of twinned pairs to in the phenomenological theory of martensite arrangement as reporter variation pairs [21-26].

Although numerous variation sets may shape from an austenite crystal every variation pair is comparable, in this way it is feasible for the austenite to change a single correspondent variation pair which accommodates the change strain by twin formation. This however does not happen in an un-one-sided (no outside pressure) test. What is watched is that austenite changes into a pretty much arbitrary orientation distribution of variant sets. An expansion mechanical imperative must be considered to inspect the driving force for the observed appropriation of variant sets. The fundamental system driving such a circulation comes from the minimization of the microscopic visible state of the bulk material [21].

The change result of the austenite is in various areas of the crystal changes in such a way, to the point that there is no microscopic shape change. This behaviour is stated as self-accommodation. Self-accommodation is a game plan of martensitic variations to such an extent that the sum of their displacements with the boundary suffers no net displacement. It is conceivable to put a self-accommodation arrangement within an austenitic matrix and not persuade any microscopic strains [27]. Self-accommodation is a main property for all thermoelastic changes as it limits interface stresses guaranteeing interface coherence and versatile accommodation of strains. In other Words, if self-accommodation did not happen as the interface progresses the pressure in the interface would proceed to rise and rapidly

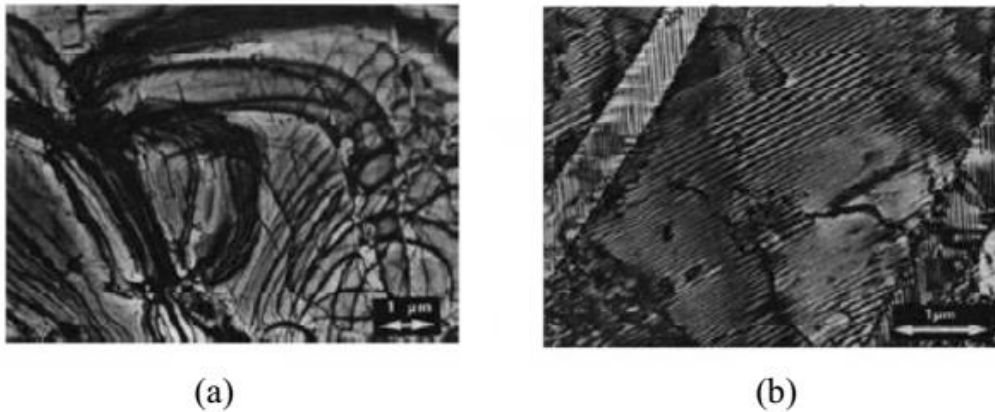
outperform its yield quality bringing about plastic deformation and lost coherency. Such is the situation in twinned ferritic martensites, where despite that the twinning procedure is reversible, the arrangement of twins in with the end goal that the interface pressures result in plastic deformation and therefore a non-recoverable change.



**Figure 1.5** Two-dimensional lattice schematic of monoclinic structures [28].

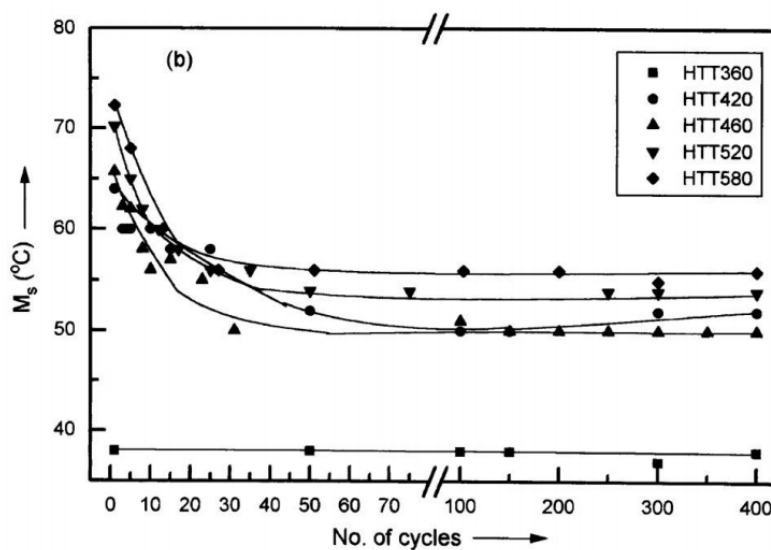
In shape memory alloys the thermoelastic change outcomes in a twinned structure. Figure 1.5 demonstrates a simplified depiction of two equal monoclinic variations isolated by a twin boundary. Although more strain could be suited by the extra interpretation related with the arrangement of an incoherent twin the interfacial energy of a coherent twin boundary is by and large a order of magnitude lower than the energy of an incoherent twin subsequently extra energy is required to shape incoherent twins [22]. Subsequently the arrangement of coherent twin interfaces between martensite variants is thermodynamically positive and extra energy is required to shape incoherent twins.

Figure 1.6 are TEM micrographs of a paired NiTi shape memory alloy [2]. This Figure demonstrates the contrasts between a twinless martensitic structure Figure 1.6(a) and a finely twinned structure Figure 1.6(b). Both structures are monoclinic contrasting just by the nearness of a fine distribution of disformation twins that shape during the stage change. The detwinned structure appeared in micrograph Figure 1.6(a) resulted because of distortion by detwinning of the martensite. Another crucial part of shape memory alloy is deformation of the martensitic structure through detwinning [29].



**Figure 1.6** TEM micrographs of twinned and untwined monoclinic martensite [28].

An entirely reversible basic change requires that the parent-martensite interface be gissile in the forward and reverse directions. In addition, thermal hysteresis must be small. Ling and Owen have demonstrated that sessile dislocation loops and other defects in the matrix facilitate the movement of the interface [23].



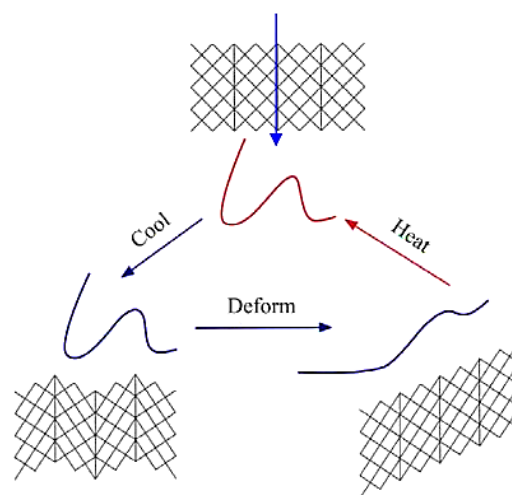
**Figure 1.7** Effects of thermal cycling through the hysteresis on the transformation temperatures of several NiTi based shape memory alloys [24].

Moreover these sessile defects increase the plastic flow pressure of the matrix hence making the accommodation of strain by slip more difficult. Mechanically the matrix is adequately fortified and the energy required to move the parent martensite interface is dropped. This has been related to NiTi bases SMAs and additionally other SMA systems. The thickness of these deformities increases with thermal cycling up as far possible resulting in a change temperatures

with expanded thermal cycles. Decreasing change temperatures with thermal cycling has been observed experimentally as it appeared in Figure 1.7 [24].

### 1.2.3 SMA Mechanical Behavior

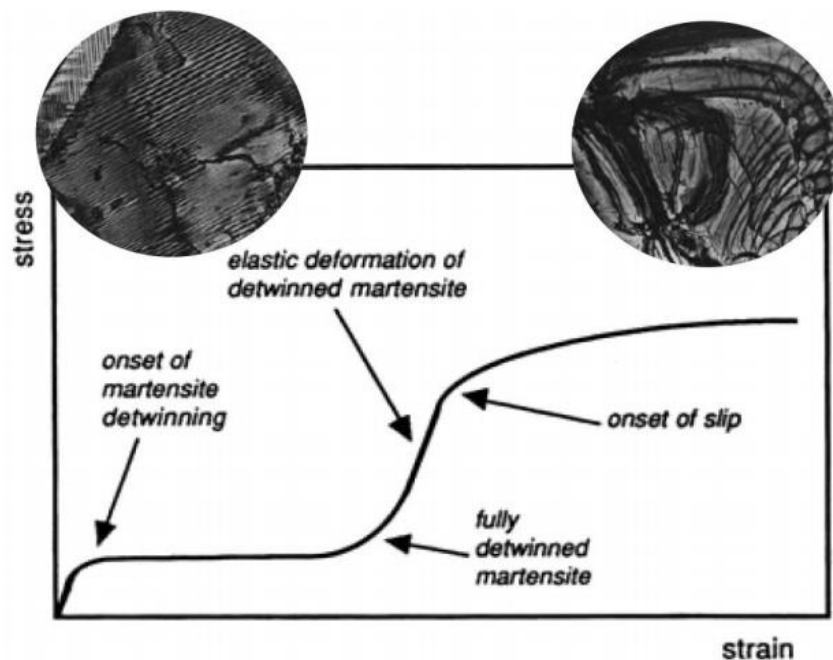
Metals that show a thermal shape memory impact deform through twin boundary movement. Recoverable deformation of the martensite by twinning responses must happen at pressure lower than those for non-crystallographically reversible responses. Non-crystallographically reversible reactions incorporate dislocation generation and movement. Basically, twins are shaped in the martensite during the forward response isolating comparable variants. Numerous martensite variants arrangements is driven by the minimization of the net transformational stresses [30]. Consequently, twins are available in the microstructure after change; hence, nucleation by a connected shear is not necessary in contrast to standard deformation twinning. Deformation happens by the development of variants most positively lined up with the biggest principle shear part of Schmidt factor at the cost of those with the lowest element. This mechanical is generally referred to as detwinning. Deformation by this system decreases the quantity of twins in the alloys [31]. This is schematically appeared in Figure 1.8. The arrow shows twin planes in the schematic. An allied shear outcomes in the twin boundary movement toward the ordinary path to the shear and the development of a relating variant. Upon completely detwinning the alloy, the material theoretically exists as a solitary variant despite the fact that the degree of detwinning relies upon crystal structure and the related number of proportional variants and additionally the current variants circulation preceding deformation [32].



**Figure 1.8** Deformation and shape recovery by detwinning (twins marked with arrows) [21].

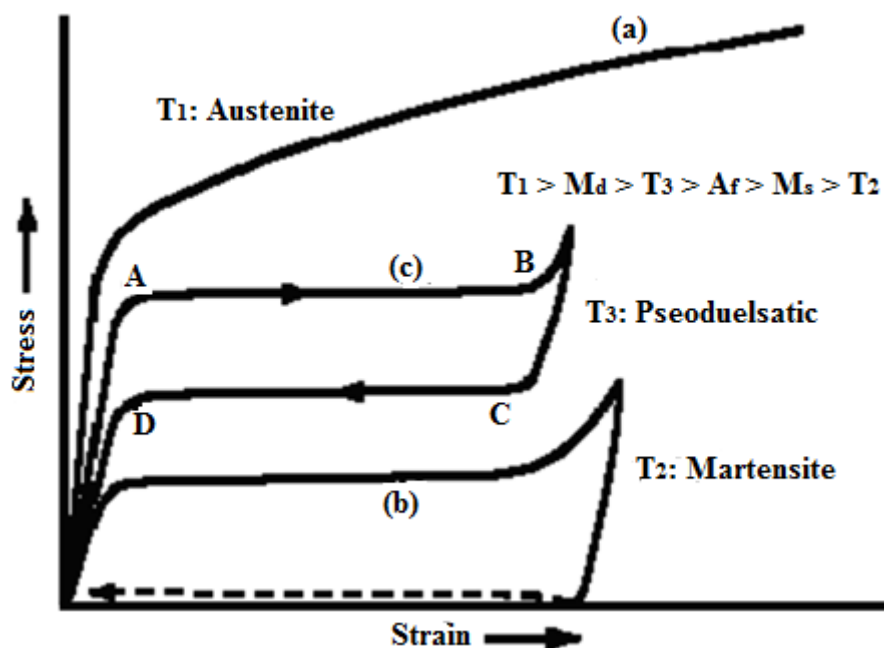
The distinctive macroscopic mechanical behavior of shape memory alloy is shown in Figure 1.9. This figure is a schematic of the universal pressure strain bends show in these systems below and the  $M_f$  temperature. It should be observed that the alloy's composition, thermal and mechanical history may change this curve. It is too conceivable to have different dynamic deformation systems, which will influence the work solidifying rate during the detwinning area of this curve. This figure shows the perfect case for the shape memory effect [33].

The early part of the pressure strain curve is credited to flexible deformation of the unreformed martensite. After achieving basic pressure, detwinning of the martensite starts. The detwinning stress is free of twin thickness and in this way an area in the stretch strain curve exists in which the stress required to twist the material is autonomous of strain [34]. A critical level is come to at the point where positively arranged variants are most pervasive in the microstructure and in this way reactants of the detwinning responses that supply the growth of the positively arranged twins are devoured. At this point twins with comparable Schmidt influence may block on each other. The outcome is an increase in stress-strain relationship that is ascribed to flexibly deforming the detwinned martensite. A moment yield point is apparent at which the critical stress for slip is come to. Non-reversible deformation mechanisms are dynamic in this area therefore the strains are not recoverable by shape memory processes [35].



**Figure 1.9** Isothermal stress strain behavior of a typical SMA strained in the fully martensitic state [3].

Figure 1.10 is a progression of hypothetical stress-strain curve that graphically show to the three-particular deformation behavior exemplified by SMAs and the temperatures at which they might be active [25]. At temperature over the  $A_f$  the Shape memory alloy is austenitic and deformation happens by elastic loading of the austenite taken after by slip. Beneath the  $M_f$  temperature the shape memory alloy is completely martensitic and deformation happens by the detwinning systems portrayed previously. SMAs exhibit an unprecedented super elastic impact which happens when the material is deformed over the  $M_d$  temperature and under the  $M_d$  temperature (Figure 1.10). The  $M_d$  temperature is characterized as the temperature at which mechanical stresses can instigate a martensitic change. Ensuing evacuation of the external stress brings about a non-diffusional inversion to the thermodynamically stable parent stage. Versatile strains achievable in NiTi Shape memory alloys are 20X those of carbon spring steels [36].



**Figure 1.10** Stress strain behavior showing the three main deformation regimes active in SMAs [37].

#### 1.2.4 Thermoelastic Shear Transformations

The thermodynamics of shape memory alloys and the significant shear changes is an unpredictable subject, which includes a composition between the chemical and non-chemical driving force. It is expressed that martensite shapes from the parent stage by an absolutely diffusionless shear change. The change front progresses by shear nuclear movements and the

interface between the martensite and parent stage is lucid [38]. Basically, the martensite brings about a net shape change of every identical variant. As on account of NiTi addressed previously a cubic structure changes to a monoclinic or orthorhombic. This net shape change results in an accommodation strain. The local strain around every variant can be accommodated plastically, flexibly or as a blend of both. This marvel has been assessed and the key thermodynamic parameter distinguished by Reed and Abbaschian [26].

Table 1 was arranged from production [1, 22, 39]. The listed features are basic to SMA systems. The mechanical and structural features already have been expressed. Therefore, as a result of these structural features a particular set of thermodynamic characteristics arises.

**Table 1.1:** Some features common to metallic thermal SMAs.

<b>Structural features of SMAs [1].</b>
<ul style="list-style-type: none"> <li>• Ordered parent &gt; ordered martensite</li> <li>• Martensitic change is thermoelastic</li> <li>• Martensite is crystallographically reversible</li> </ul>
<b>Thermodynamic features characterized by Dunne and Wayman [40]</b>
<ul style="list-style-type: none"> <li>• Small chemical driving force at Ms</li> <li>• Small transformational volume and shape change</li> <li>• High flow stress parent matrix</li> </ul>
<b>Further mechanical features characterized by Ling and Owen [22]</b>
<ul style="list-style-type: none"> <li>• Parent –martensite interface must be glassily in both transformation directions</li> <li>• Pre-martensite elastic softening</li> </ul>

A key condition for the shape memory effect is that the change must happen reversibly. Accommodation of the transformational strains contiguous the interface could be elastic, plastic, or a blend of both. For the situation where the lion’s share the strain is suited flexibly the interface can move in both directions referred to as a thermoelastic change. Thus, the



chemical driving force need to drive the response is small. This is the situation in alloys showing the shape memory effect, which are all the more accurately characterized as thermoelastic changes [41].

In thermoplastic changes, the change strain is suited plastically because of a low flow stress in the parent stage and an extensive transformation strain. For this situation, the changes on the forward and reverse happen at higher chemical driving forces through the nucleated and quick development of the martensite. Normally, individual shear plates are nucleated at imperfections and grow irreversibly [42].

Consequent thermal cycles result in new plates nucleating as opposed to reversible interface movement. A great example of such a situation in the martensitic change of carbon steels where the thermal hysteresis is large and the vast majority of the transformational strains are accommodated plastically [43].

Ortin and planes gracefully treated thermodynamics of thermoelastic effects and systematically characterized conditions for a thermoelastic energy balance. Thermoelastic changes are restricted by the chemical free energy. At equilibrium, it would be predictable that the change happens when the chemical free energy of the parent stage is a minor amount larger than that of shear stage [44]. Though it's not the case that has been shown. In reality, the chemical driving force in thermoelastic change are opposed by non-chemical forces are nearly equal. Following the approach and notation exhibited by Ortin and planes, thermodynamic balance is shown by the following equation.

$$\Delta G_{p..m} = -\Delta G_{ch} + \Delta G_{nch} = 0 \quad (1.2)$$

$G_{p..m}$  is the molar free energy of change,  $G_{ch}$  is the molar chemical free energy and  $G_{nch}$  is a moral non-chemical energy. At equilibrium, the molar free energy of change is equal to zero therefore the chemical contributions are equal to the non-chemical contribution [45].

The non-chemical contributions comprise of few components, the most conspicuous are the elastic free energy and work done against frictional forces. The frictional and elastic terms comprise of a few parts [46].

The elastic and frictional terms comprise of a numerous part. It was expressed that a condition for thermoelastic change is that most of the transformational strain are accommodated for

elastically. So, there exists an interfacial energy related with the parent martensite interface and also the interfaces between variants (twin boundaries) [47].

$$\left. \begin{aligned} \Delta G_{nch} &= \Delta G_{el} + E_{friction} \\ \Delta G_{ch} &= \Delta G_{nch} \\ \Delta G_{ch} &= \Delta G_{el} + E_{friction} \end{aligned} \right\} \quad (1.3)$$

Both contributions are reversible along these lines they have been assembled into the elastic term despite the fact that the interfacial energy is not genuinely an elastic contribution. This is in accordance with the convention set by Ortin and Planes [48].

Frictional energy losses are non-reversible losses primarily because of interface movement. This term might be dealt with as irreversible work done on the system. Three noteworthy parts have been recognized by Olsen and Cohen [37, 49]. These incorporate

1. frictional stresses need to move interfaces
2. irreversible free energy identified to defects made during the change and
3. frictional stresses need to move interface.

It is significant to note that assuming all or the greater part of the accommodation happens plastically the flexible term will be small and the frictional energy losses term will be the primary restricting energy to the chemical driving force. Thus, a huge non-reversible hysteresis will be obvious just like the case in carbon steels as stated before [50].

The fundamental thermodynamics and mechanism which control general conduct of shape memory alloys have been observe. Essentially the structural relationships between the parent in shear stages result in thermomechanical characteristics special to shape memory alloys. Also, driving forces and the relating thermodynamic for thermoelastic changes have been briefly explained and contrasted with thermoplastic changes [51]. A persistent issue that will be explained in consequent chapters brings up that plastic accommodation by non-recoverable procedures hinder the sought-after attributes of shape memory alloys, basically the materials capacity to do work. Principally, under specific materials where a huge part of the change strain is accommodated plastically. The mix of structure, thermodynamic and mechanical properties are utilized as devices to comprehend deformation and propose conceivable deformation mechanical while recognize possible target area for future alloy improvement [52].

### **1.3 Organization of Thesis**

Chapter 1 includes the introduction and the properties of shape memory materials.

Chapter 2 includes the theoretical frameworks and literature review on the research topic.

Chapter 3 includes the simulations and analyses run to visualize the effect of forces on spring.

Chapter 4 includes the results and discussion.

Chapter 5 includes the conclusion and future work.

## CHAPTER 2

### THEORETICAL FRAMEWORK

#### 2.1 Design of a Spring

Spring is an important machine element that comes in action whenever rigidity and small deflections are required provided that it does not undermine the functionality of the product. Flexibility is usually required when internal machine parts needs to be controlled. These flexibility is attained by store and release of energy on application of force or torque on the spring [53].

Usually springs are classified as follow:

1. Wire springs
2. Flat springs
3. Special-shaped springs

The first type of spring include helical springs of round and square wires which can withstand under compressive, tensile or compressional loads. Wound motor- or clock-type power springs, Cantilever and elliptical springs, flat spring washers are the types of flat springs also known as Belleville springs [54].

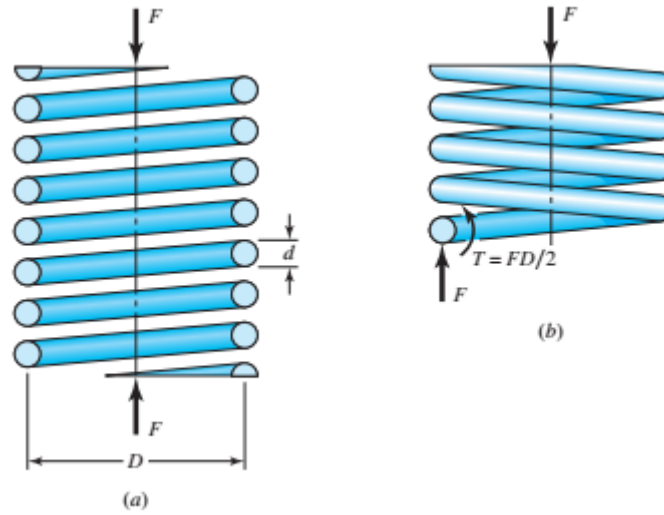
#### 2.2 Stresses in a Spring

The process of evaluation of stress in spring begins with defining parameter. Let  $d$  and  $D$  are the *wire diameter* and *mean coil diameter* respectively. A spring is shown in Figure 2.1(a) is loaded with axial force  $F$ . A section of spring in Figure 2.1(b) is taken from it and in the state of equilibrium, the shear force acting on it is  $F$  and torsional moment acting is  $T = FD/2$  .

The shear stress that could appear in this situation is evaluated by adding the shear force and the torsional force as follow:

$$\tau_{\max} = \frac{Tr}{J} + \frac{F}{A} \quad (2.1)$$

This shear stress appears at the inner side of the spring.



**Figure 2.1** (a) Spring with axial load; (b) free-body diagram of spring under a direct shear and a torsional shear [54].

Doing the following suitable substitutions in Eq. (2.1):

$$\tau_{\max} = \tau \quad T = FD/2 \quad r = d/2 \quad J = \pi d^4/32 \quad A = \pi d^2/4 \quad (2.2)$$

The Eq. (2.1) becomes:

$$\tau = \frac{8FD}{\pi d^3} + \frac{4F}{\pi d^2} \quad (2.3)$$

Assuming the *spring index* as

$$C = \frac{D}{d} \quad (2.4)$$

It is the computation of coil curvature. Using Eq. (2.3) and Eq. (2.4) together yields:

$$\tau = K_s \frac{8FD}{\pi d^3} \quad (2.5)$$

where  $K_s$  is shear–correction factor and is as follow:

$$K_s = \frac{2C+1}{2C} = 1 + \frac{1}{2C} \quad (2.6)$$

### 2.3 Curvature Effect of a Spring

The Eq. (2.5) is used when the wire is straight which not the case is in helical spring hence, curvature of the wire is the reason of increase in stress on the inner side of the spring coil. It is applicable in the same way as stress correction factor. In case of static loading, this factor is meaningless since it causes strain hardening while in dynamic loading, which usually the normal function of spring is, the fatigue needs it to be included in the stress evaluation [55].

It is important to find curvature factor for the section of the spring but it is usually found through direct shear stress. Assuming  $K_s$  in Eq. (2.5) is replaced by  $K$  which can solve both issues of curvature and direct shear then it can be defined as

$$K_w = \frac{4C-1}{4C-4} + \frac{0.615}{C} \quad (2.7)$$

$$K_B = \frac{4C+2}{4C-3} \quad (2.8)$$

The factor represented by Eq. (2.7) is known as *Wahl factor* and the factor represented by Eq. (2.8) is known as *Bergsträsser factor*. Both factors differ by 1% hence Eq. (2.8) is better to use. The true curvature factor is obtained by cutting out the effect of direct shear stress [56]. By using Eq. (2.8) with Eq. (2.6), the curvature correction factor is given by

$$K_c = \frac{K_B}{K_s} = \frac{2C(4C+2)}{(4C-3)(2C+1)} \quad (2.9)$$

To include all stress-correction factors, simply multiply these with Eq. (2.5)

$$\tau = K_B \frac{8FD}{\pi d^3} \quad (2.10)$$

This is the maximum shear stress that appears on the spring.

### 2.4 Deflection in a Spring

The relation of deflection and force can be obtained by *Castigliano's theorem*. The total strain energy of a spring consists of summation of shear and torsional components. Hence the strain energy is

$$U = \frac{T^2 l}{2GJ} + \frac{F^2 l}{2AG} \quad (2.11)$$

Replacing the following substitutions:

$$T = FD/2 \quad l = \pi DN \quad J = \pi d^4/32 \quad A = \pi d^2/4 \quad (2.12)$$

The result of strain energy Eq. (2.11) becomes:

$$U = \frac{4F^2 D^3 N}{d^4 G} + \frac{2F^2 DN}{d^2 G} \quad (2.13)$$

In the above relation of strain energy,  $N = N_a$  = number of active coils . Applying *Castigliano's theorem* to evaluate deflection  $y$ :

$$y = \frac{\partial U}{\partial F} = \frac{8FD^3 N}{d^4 G} + \frac{4FDN}{d^2 G} \quad (2.14)$$

Using Eq. (2.4) with Eq. (2.14)

$$y = \frac{8FD^3 N}{d^4 G} \left( 1 + \frac{1}{2C^2} \right) \approx \frac{8FD^3 N}{d^4 G} \quad (2.15)$$

The parameter *spring rate* is also known as *scale* of the spring, which is  $k = F/y$ , and can be defined as

$$k \approx \frac{d^4 G}{8D^3 N} \quad (2.16)$$

## 2.5 Stability of a Spring

When the axial force is applied on a structure, it buckles. Same is the case with the spring, if the axial force is too large, the spring will buckle [57]. To avoid this, critical deflection  $y_{cr}$  is given as:

$$y_{cr} = L_0 C_1' \left( 1 - \sqrt{1 - \frac{C_2'}{\lambda_{eff}^2}} \right) \quad (2.17)$$

Samónov [58] stated that Eq. (2.17) was used by Wahl [59] and the same equation is experimentally verified by Haringx [60, 61]. The factor  $\lambda_{\text{eff}}$  in Eq. (2.17) is known as *effective slenderness ratio* and is specified by:

$$\lambda_{\text{eff}} = \frac{\alpha L_0}{D} \quad (2.18)$$

Whereas the factors  $C'_1$  and  $C'_2$  are dimensionless elastic constants given by:

$$C'_1 = \frac{E}{2(E - G)} \quad (2.19)$$

$$C'_2 = \frac{2\pi^2(E - G)}{2G + E} \quad (2.20)$$

The Eq. (2.18) contains a factor  $\alpha$  known as end-condition constant which depends upon the method of support on the ends of spring. The table represents the normal values for  $\alpha$

**Table 2.1:** End Condition of a Spring

End Condition	Constant $\alpha$
Support is flat parallel plate	0.5
One end supported by plane at right angle to spring axis	0.707
Both ends hinged	1
One end free; other end clamped	2

For the stability to be absolute, the term  $C'_2/\lambda_{\text{eff}}^2$  in Eq. (2.17) must be greater than 1. This condition is met if and only if

$$L_0 < \frac{\pi D}{\alpha} \sqrt{\frac{2(E - G)}{2G + E}} \quad (2.21)$$

For conventional materials (steel), this condition is as follow:



$$L_0 < 2.63 \frac{D}{\alpha} \quad (2.22)$$

For ground and squared ends which is supported by flat parallel planes,  $\alpha = 0.5$  and  $L_0 < 5.26D$ .

## 2.6 Spring Materials

Springs are usually manufactured by the following processes:

1. Hot working
2. Cold working

The process of making spring depends upon the spring index and the size of the material. If the spring index  $C = D/d < 4$  or if  $d > 1/4$  in, then wire is not generally pre-hardened. When spring wire is wound, it causes residual stresses due to bending but these are at right angle to the torsional working stress. This residual stress is often relieved by minor heat treatment [62].

There are various materials suited for spring manufacturing including alloy steels, carbon steels and corrosion-resisting steels besides nonferrous materials like bronze, beryllium, phosphorous, copper, spring brass and many nickel based alloys [63].

The material of springs are compared by tensile strength but this differs with length of spring wire and hence true strength cannot be known unless the diameter of wire  $d$  is known. To capture this problem, a plot of tensile strength against wire diameter is drawn which is approximately a straight line when drawn on log-log graph [64]. The equation of the relation between two variables is

$$S_{ut} = \frac{A}{d^m} \quad (2.23)$$

Eq. (2.23) has  $A$  as an intercept and  $m$  as slope. It gives the least tensile strength of the spring material under consideration. Another factor, torsional yield strength is also required to estimate the strength of the spring material but generally spring material is passed on tensile strength only. It is because, by a grave estimate, the tensile yield point is in between 60 – 90 %

## Chapter 02

of tensile strength [65]. According to *Distortion-Energy* theory, the torsional yield strength is approximated as  $S_{sy} = S_y / \sqrt{3}$ . This method gives the following result for steels:

$$0.35 \leq S_{sy} \leq 0.52 S_{ut} \quad (2.24)$$

# CHAPTER 3

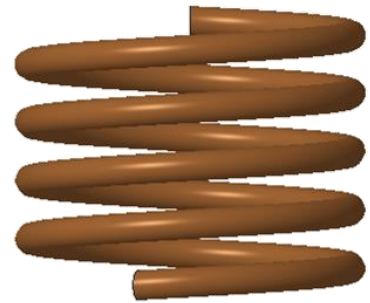
## SIMULATIONS AND ANALYSES

### 3.1 3D Modelling

The spring is designed in a commercial 3D CAD package with following parameters (Table 3.1) and CAD model is shown in Figure 3.1.

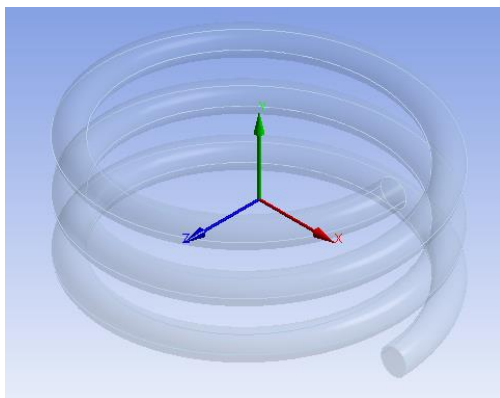
**Table 3.1 Parameters of spring**

Parameter	Value
Wire diameter, $d$ [mm]	2.16
Coil diameter, $D$ [mm]	23.1
Number of Active coils, $n$	3
Initial free length, $l_0$ [mm]	19

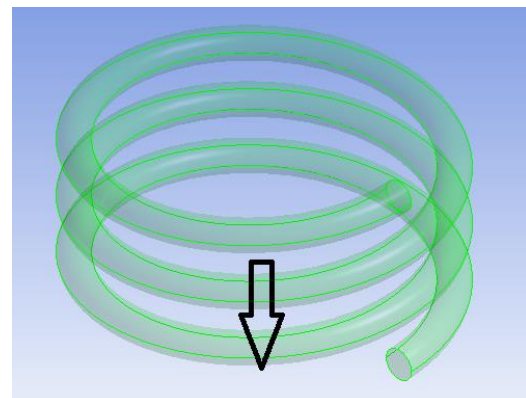


**Figure 3.1** 3D CAD model

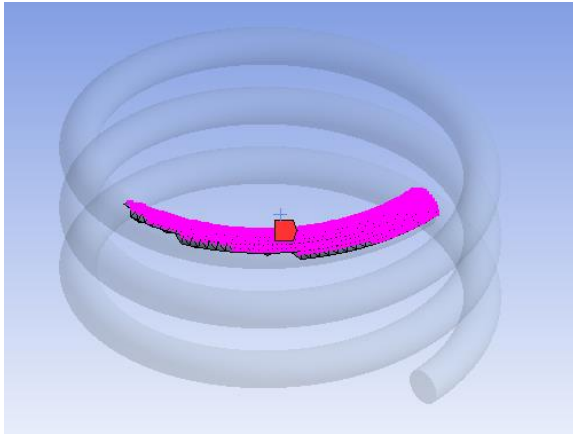
The spring is loaded with axial loading  $P$  as shown in Figure 3.3 keeping the top of the spring fixed (Figure 3.4). The rest of the spring is constrained in  $x$ - and  $z$ -directions while free in  $y$ -direction (Figure 3.5).



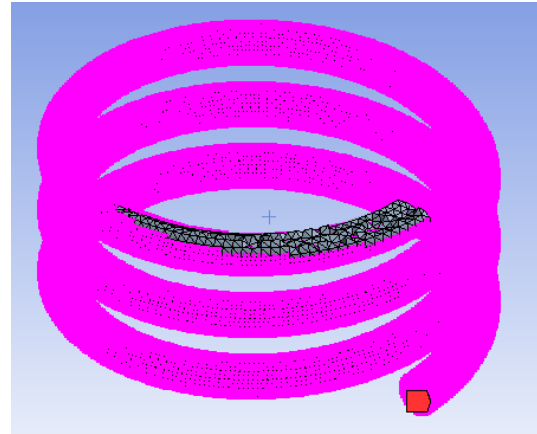
**Figure 3.2** The coordinate system for simulation



**Figure 3.3** Force  $P$  applied on lower end of spring



**Figure 3.4** Fixed nodes in all directions



**Figure 3.5** Fixed nodes in  $x$ - and  $z$ -directions and free in  $y$ -direction

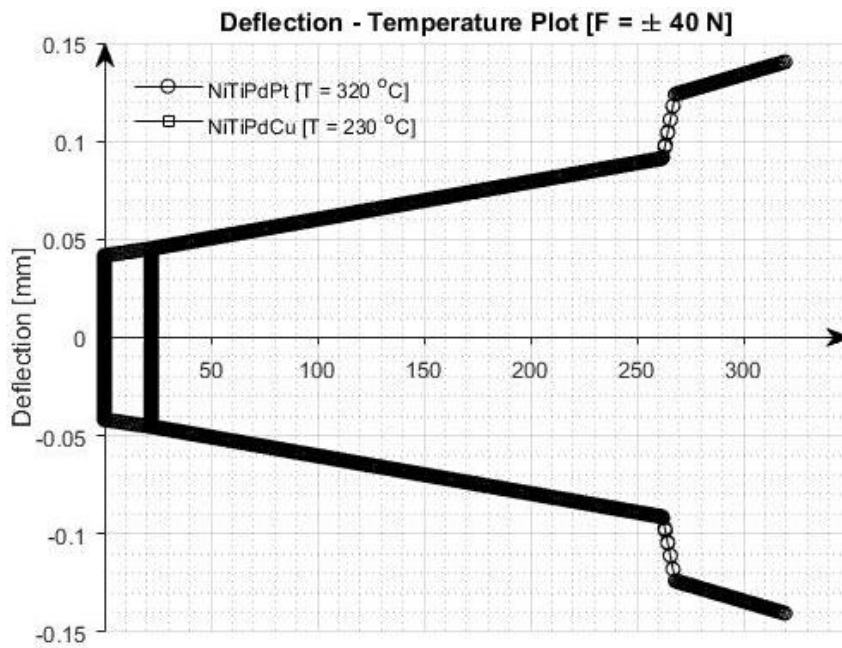
Keeping axial load applied, the temperature is increased beyond transformation temperature to introduce thermal strain. Removing temperature loading followed by removal of axial loading.

The behaviour of NiTiPdPt and NiTiPdCu has been realized by conducting experiments. These results are compared with FEA results and extrapolated. The following plots displays the relation of stress, strain, displacement, temperature and energy dissipation per unit volume.

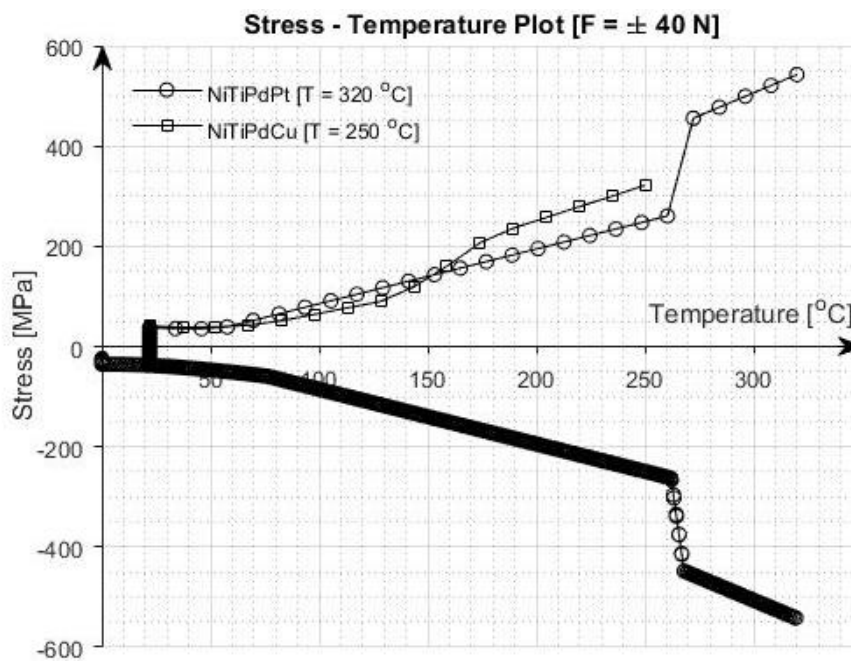
### 3.2 Simulations

#### 3.2.1 SMAs under 40 N force

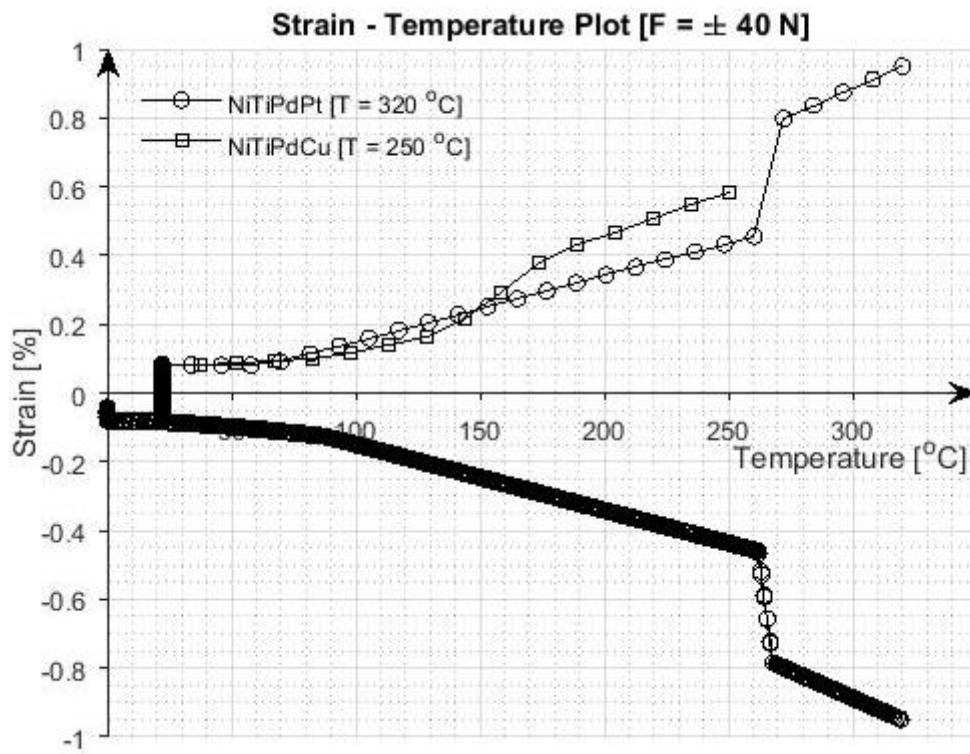
The following plots represent the results of  $\text{Ni}_{19.5}\text{Ti}_{50.5}\text{Pd}_{25}\text{Pt}_5$  at  $320\text{ }^\circ\text{C}$  and  $\text{Ni}_{15}\text{Ti}_{50}\text{Pd}_{25}\text{Cu}_{10}$  at  $250\text{ }^\circ\text{C}$  under force of  $\pm 40\text{ N}$ .



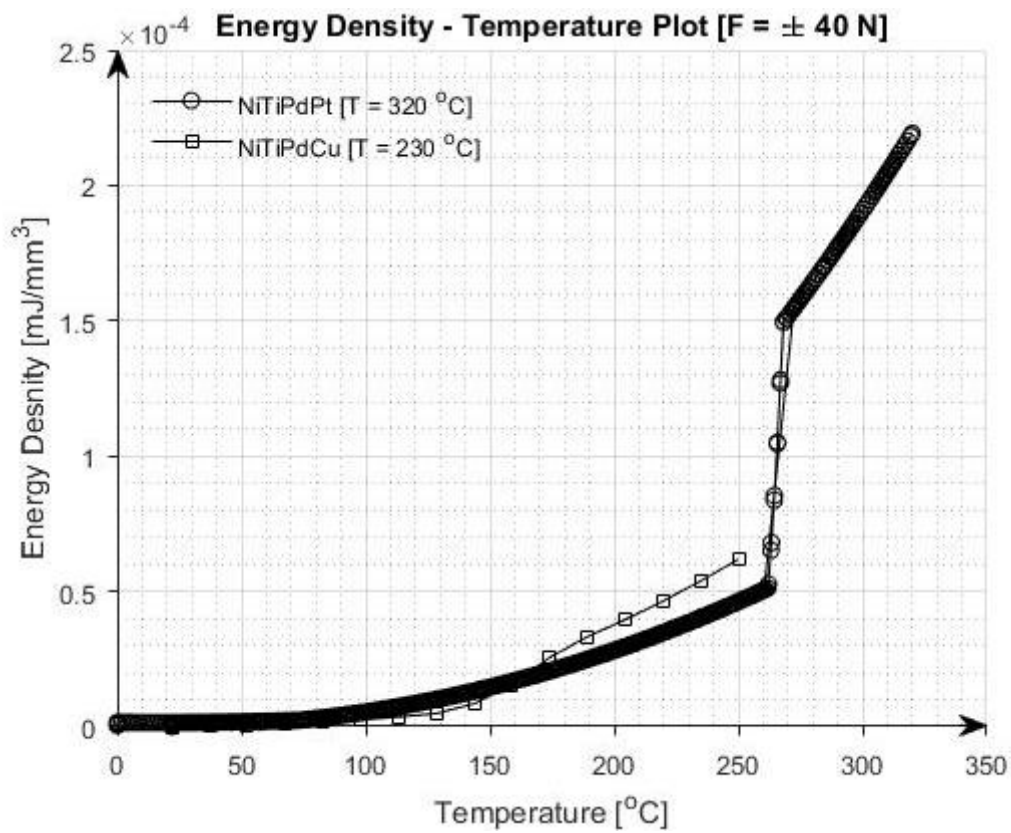
**Figure 3.6** Deflection vs. Temperature for both SMAs under  $F = \pm 40\text{ N}$ .



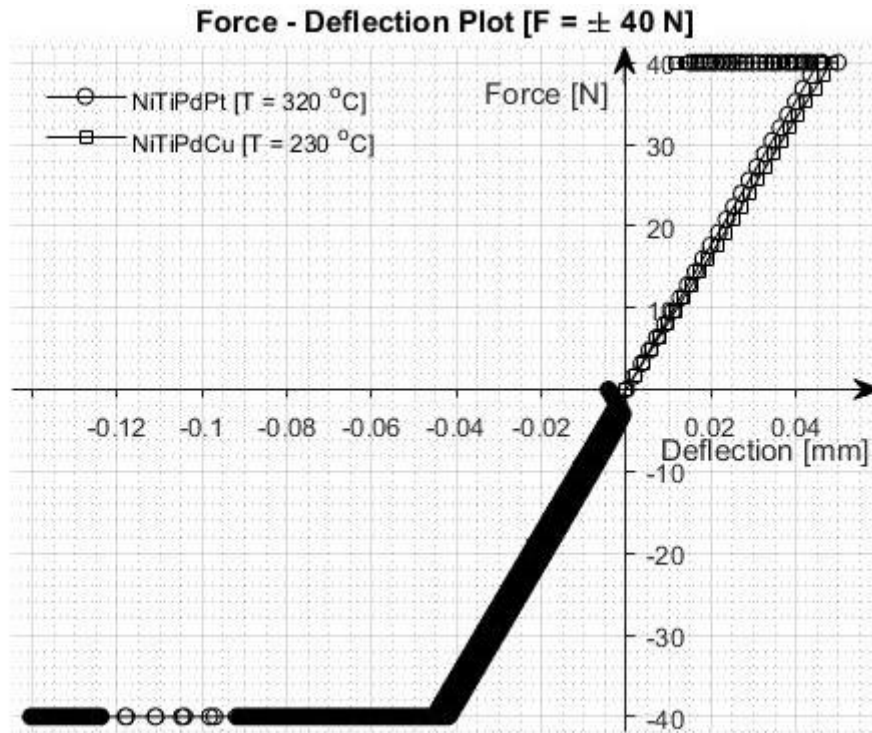
**Figure 3.7** Stress vs. Temperature for both SMAs under  $F = \pm 40\text{ N}$ .



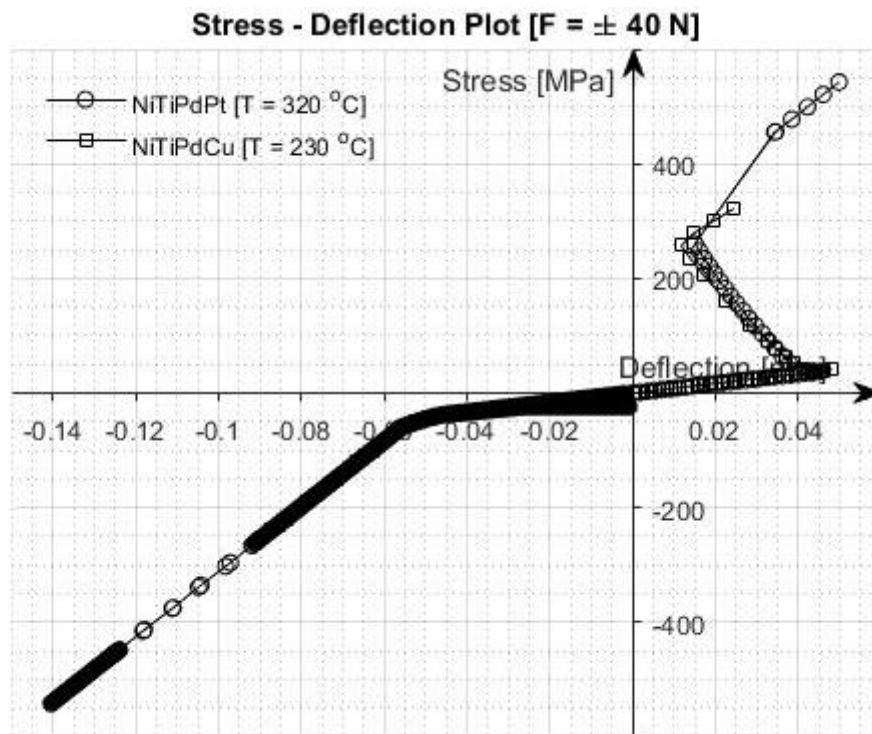
**Figure 3.8** Strain vs. Temperature for both SMAs under  $F = \pm 40$  N.



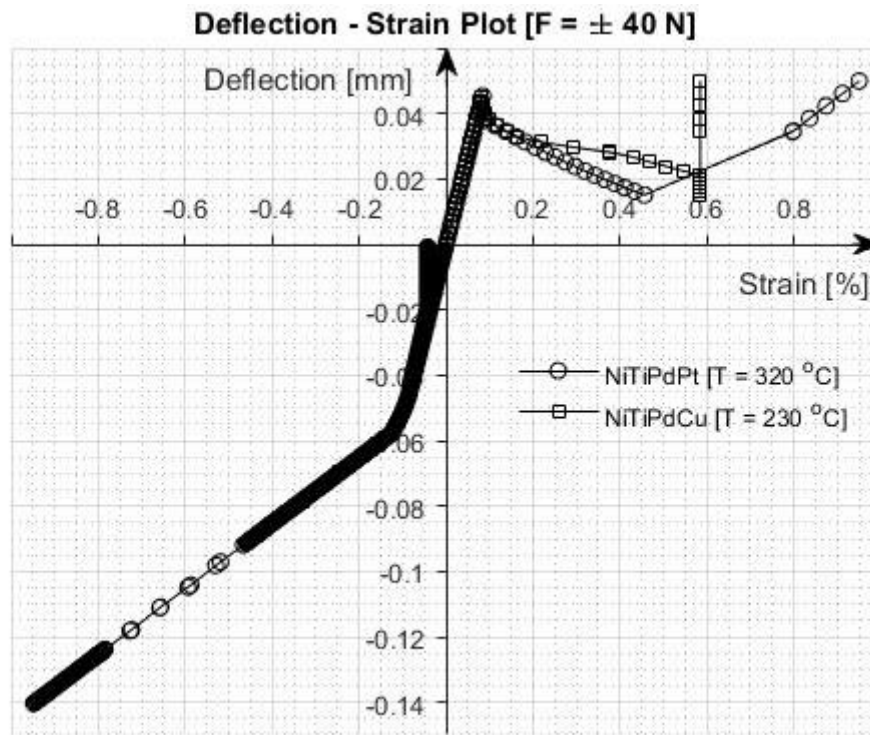
**Figure 3.9** Energy density vs. Temperature for both SMAs under  $F = \pm 40$  N.



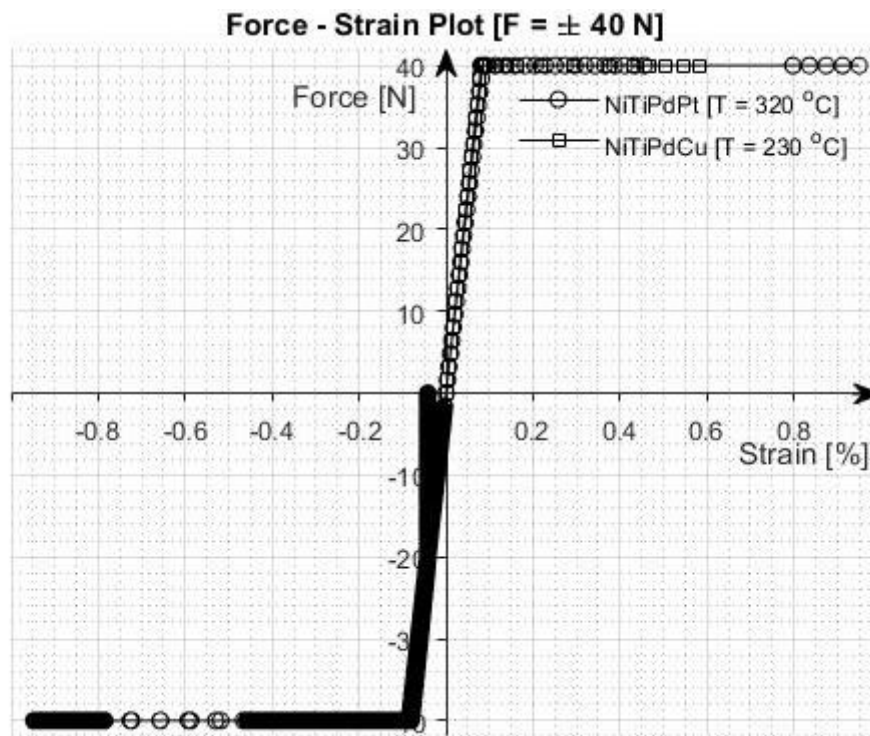
**Figure 3.10** Force vs. Deflection for both SMAs under  $F = \pm 40$  N.



**Figure 3.11** Stress vs. Deflection for both SMAs under  $F = \pm 40$  N.



**Figure 3.12** Deflection vs. Strain for both SMAs under  $F = \pm 40$  N.



**Figure 3.13** Force vs. Strain for both SMAs under  $F = \pm 40$  N.



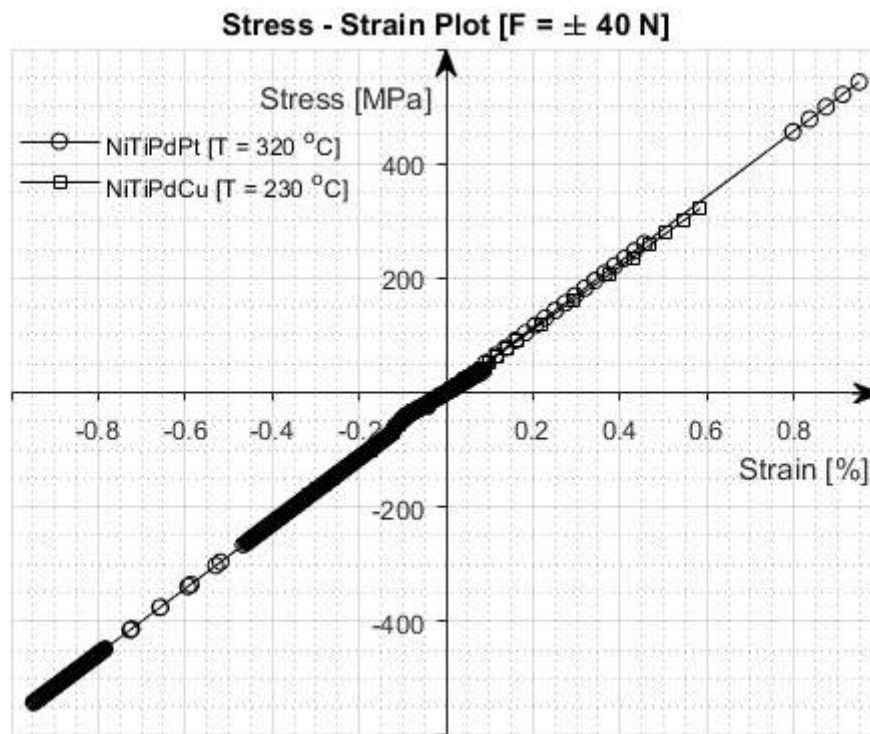


Figure 3.14 Stress vs. Strain for both SMAs under  $F = \pm 40$  N.

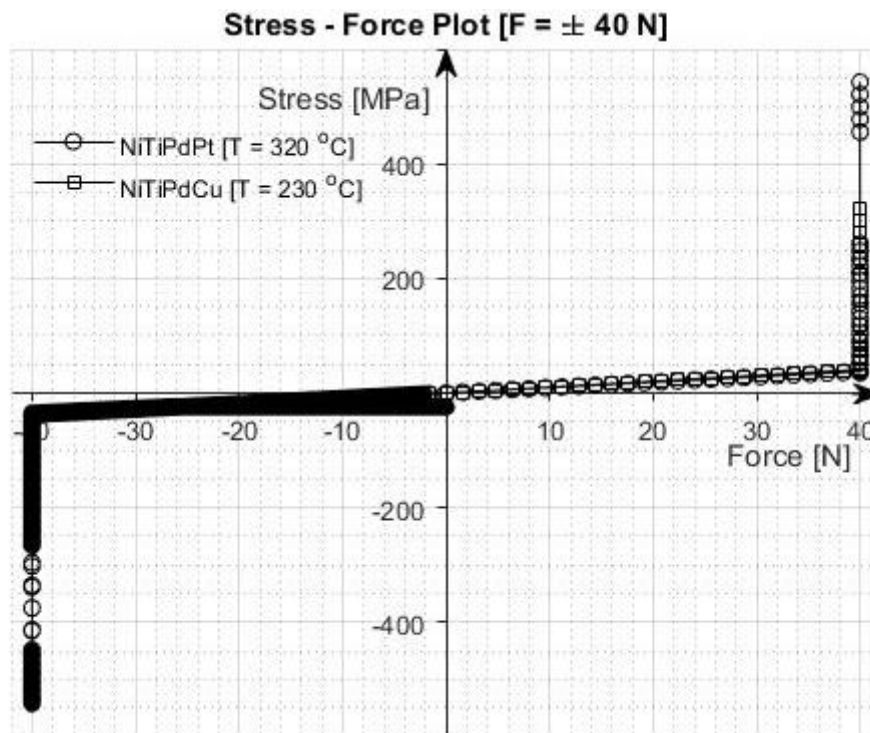
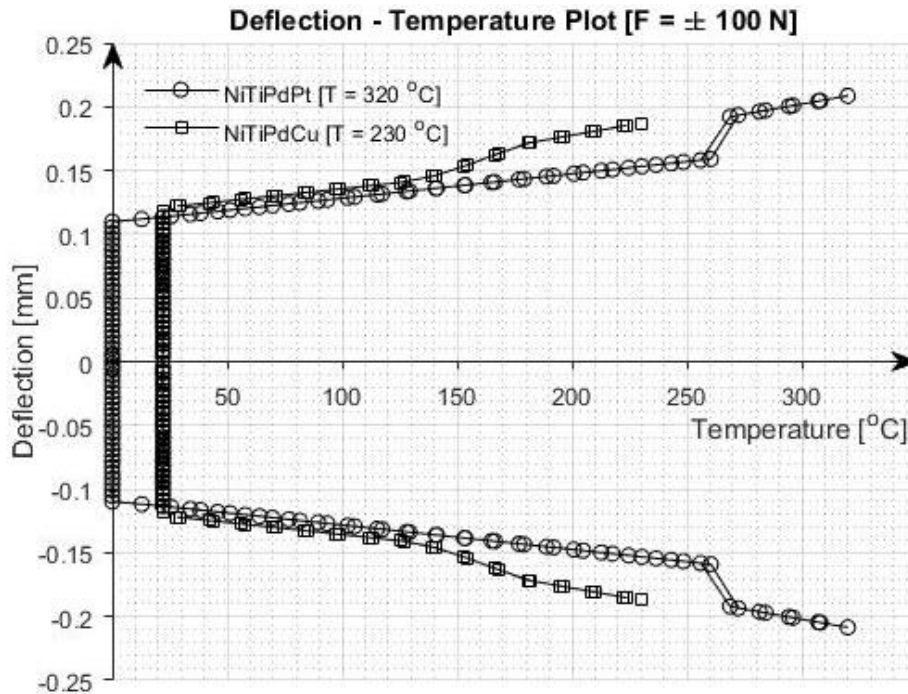


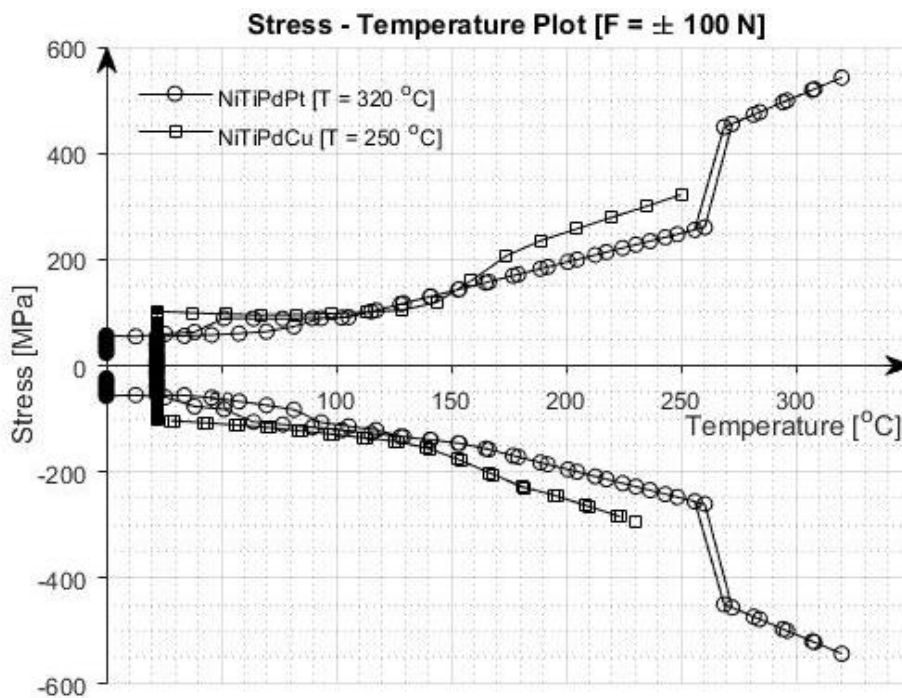
Figure 3.15 Stress vs. Force for both SMAs under  $F = \pm 40$  N.

### 3.2.2 SMAs under 100 N force

The following plots represent the results of  $\text{Ni}_{19.5}\text{Ti}_{50.5}\text{Pd}_{25}\text{Pt}_5$  at  $320\text{ }^\circ\text{C}$  and  $\text{Ni}_{15}\text{Ti}_{50}\text{Pd}_{25}\text{Cu}_{10}$  at  $250\text{ }^\circ\text{C}$  under force of  $\pm 100\text{ N}$ .



**Figure 3.16** Deflection vs. Temperature for both SMAs under  $F = \pm 100\text{ N}$ .



**Figure 3.17** Stress vs. Temperature for both SMAs under  $F = \pm 100\text{ N}$ .

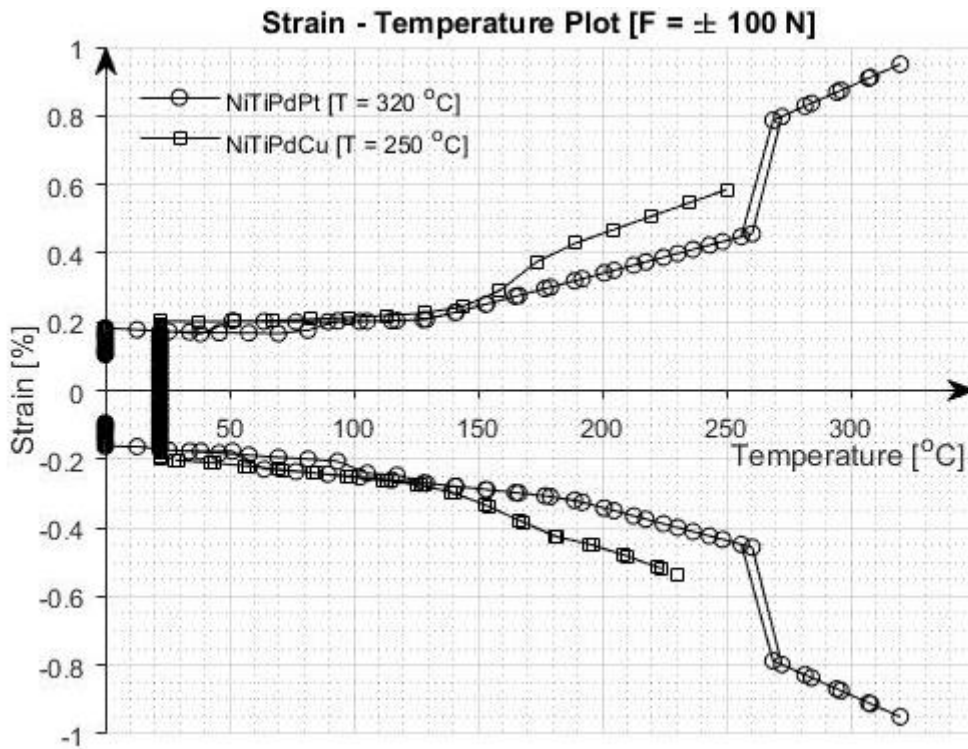


Figure 3.18 Strain vs. Temperature for both SMAs under  $F = \pm 100$  N.

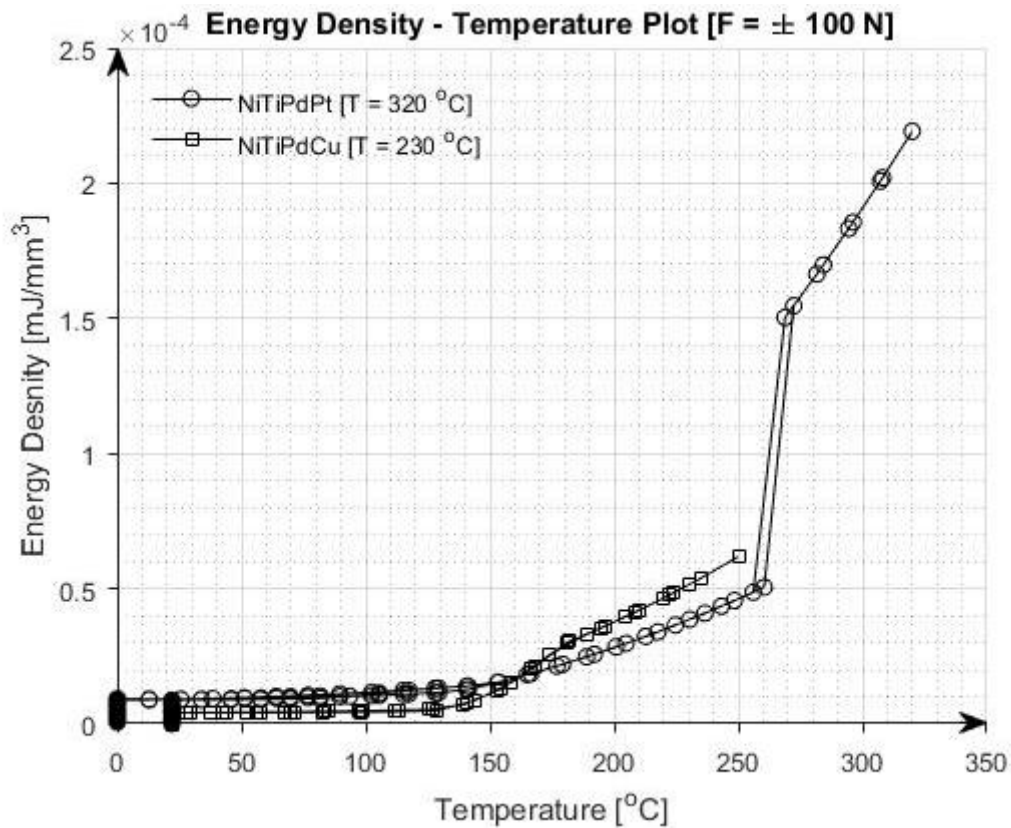
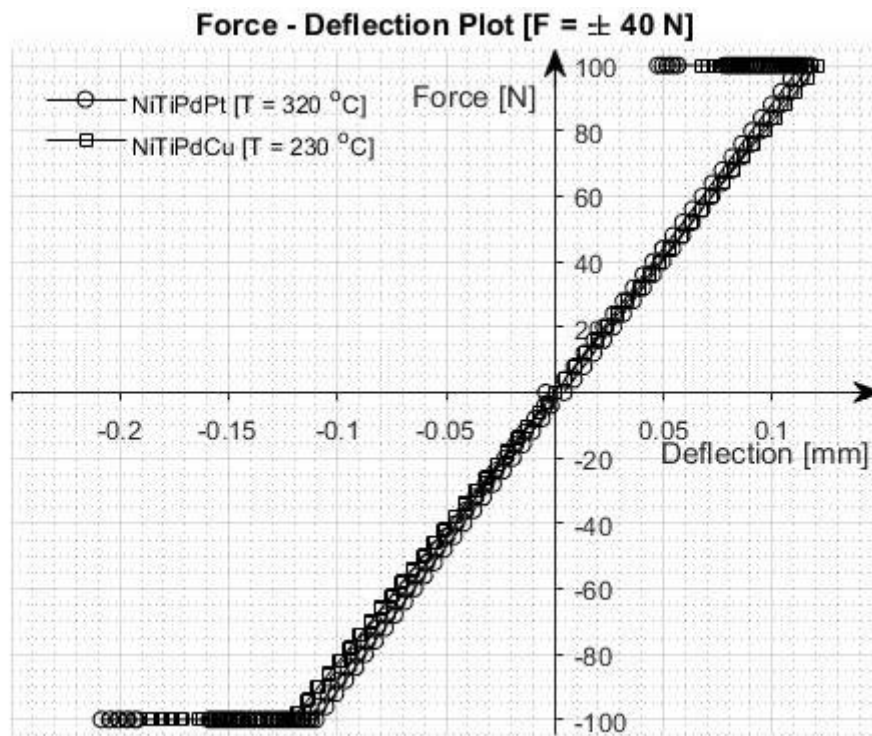
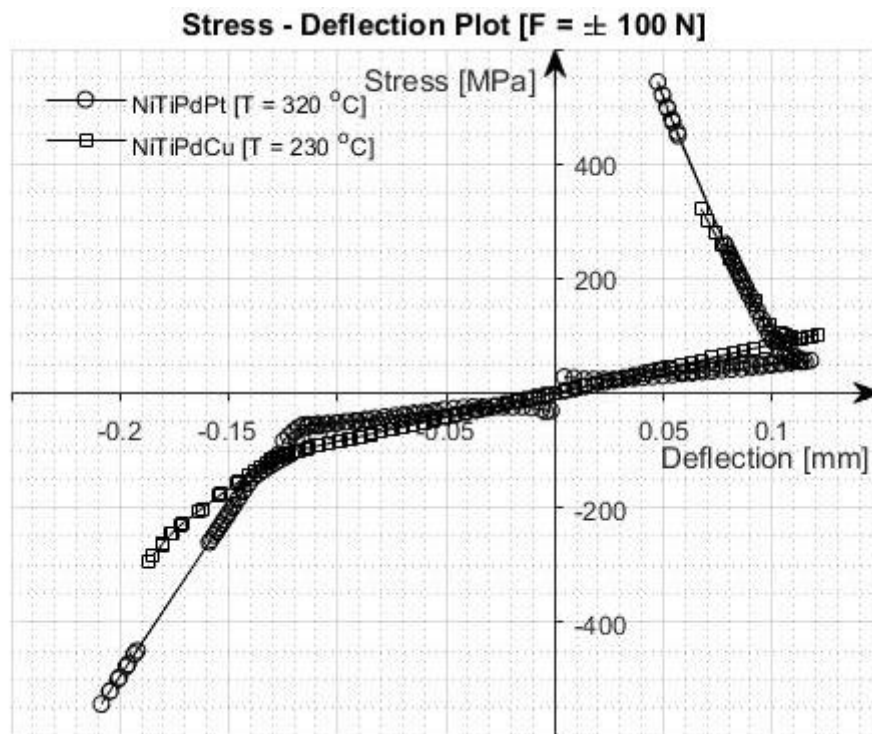


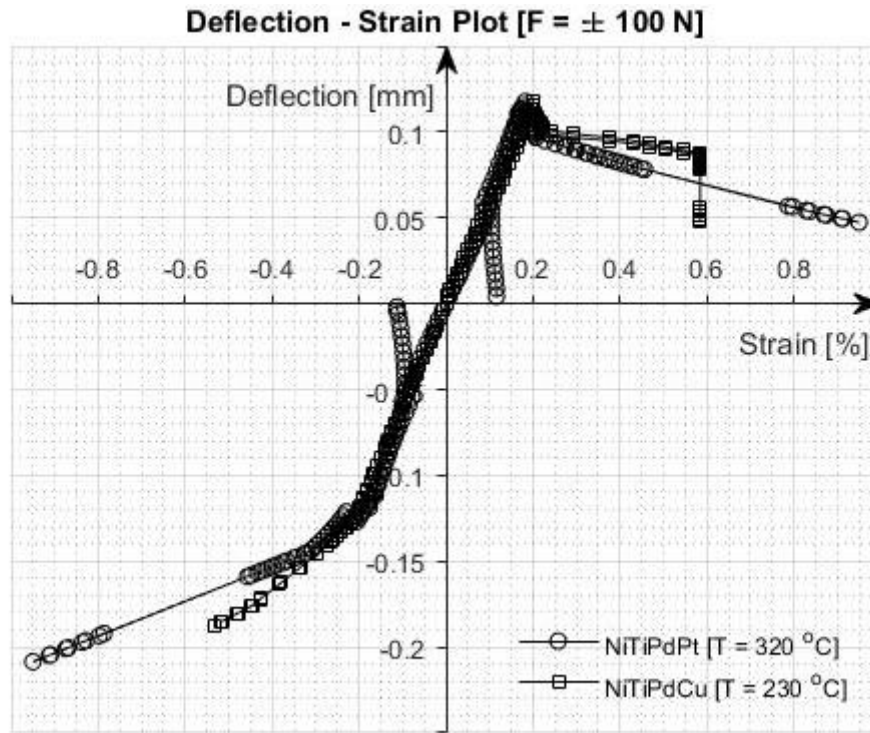
Figure 3.19 Energy density vs. Temperature for both SMAs under  $F = \pm 100$  N.



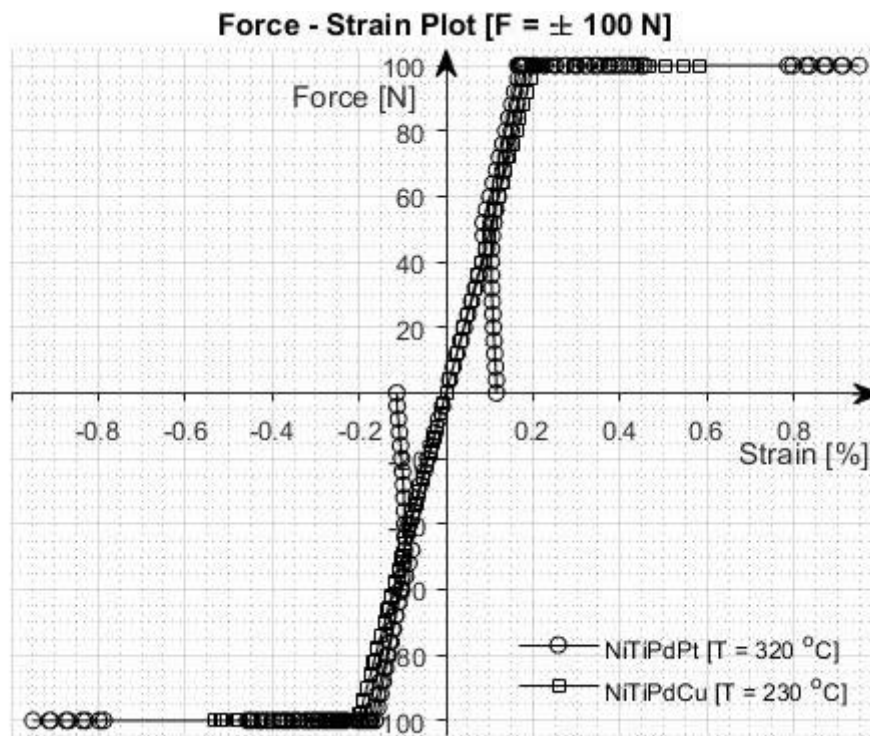
**Figure 3.20** Force vs. Deflection for both SMAs under  $F = \pm 100$  N.



**Figure 3.21** Stress vs. Deflection for both SMAs under  $F = \pm 100$  N.



**Figure 3.22** Deflection vs. Strain for both SMAs under  $F = \pm 100$  N.



**Figure 3.23** Force vs. Strain for both SMAs under  $F = \pm 100$  N.

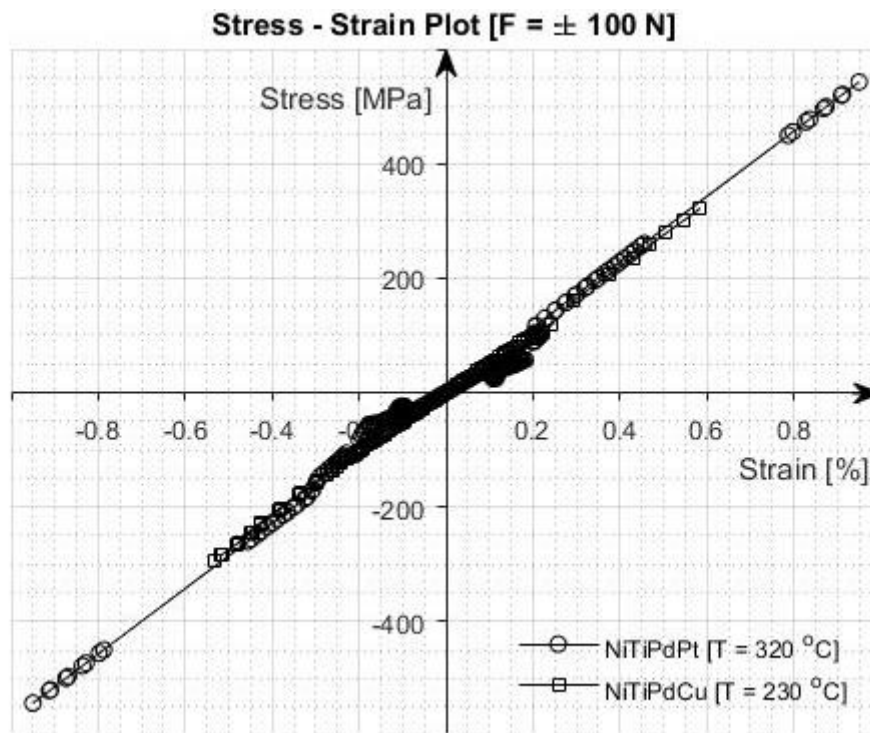


Figure 3.24 Stress vs. Strain for both SMAs under  $F = \pm 100$  N.

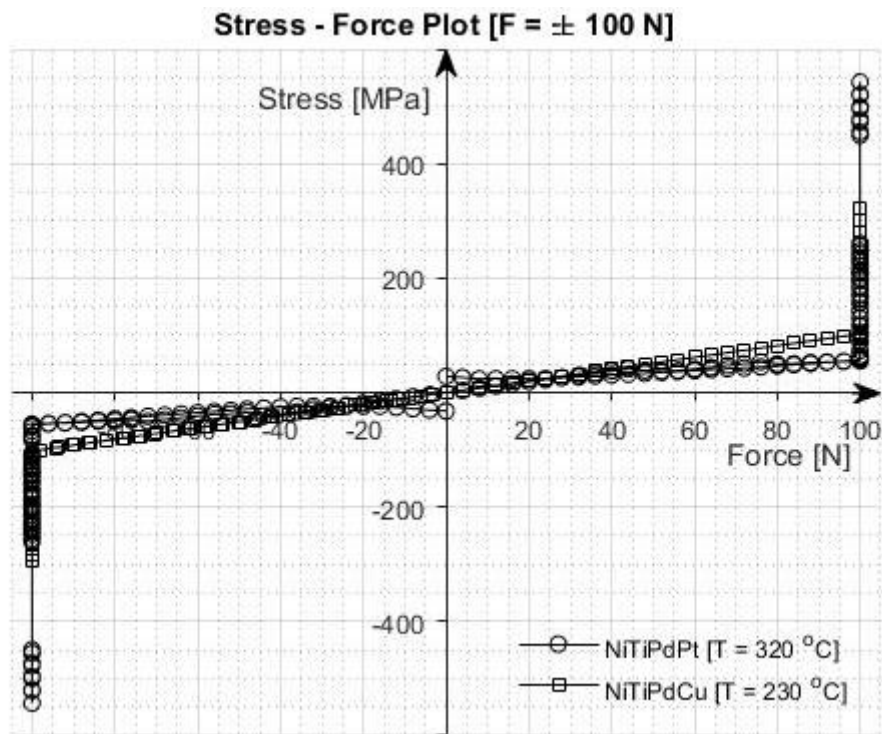
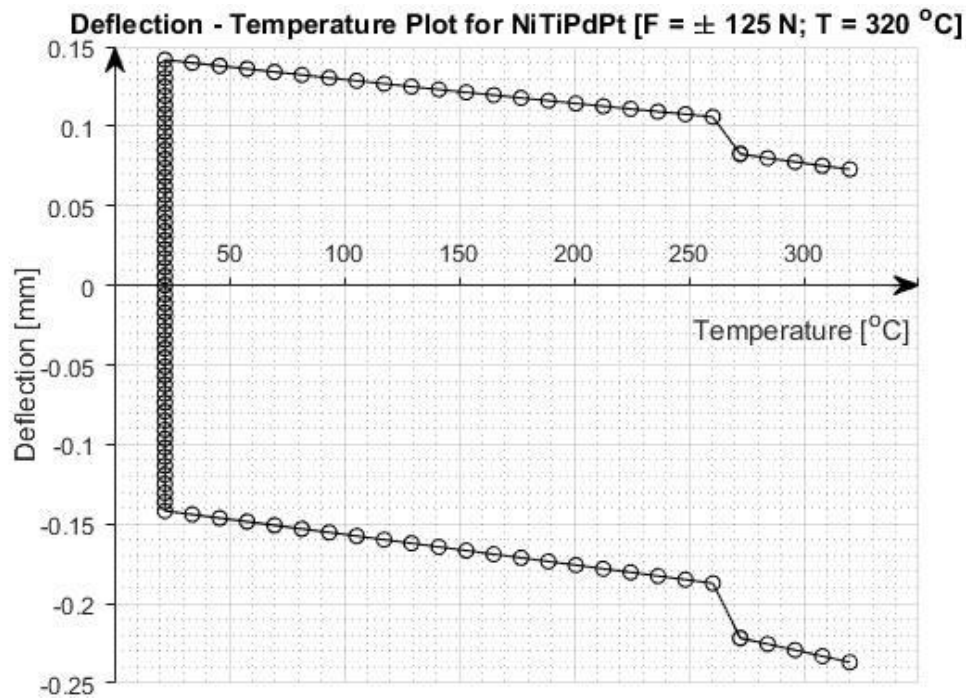


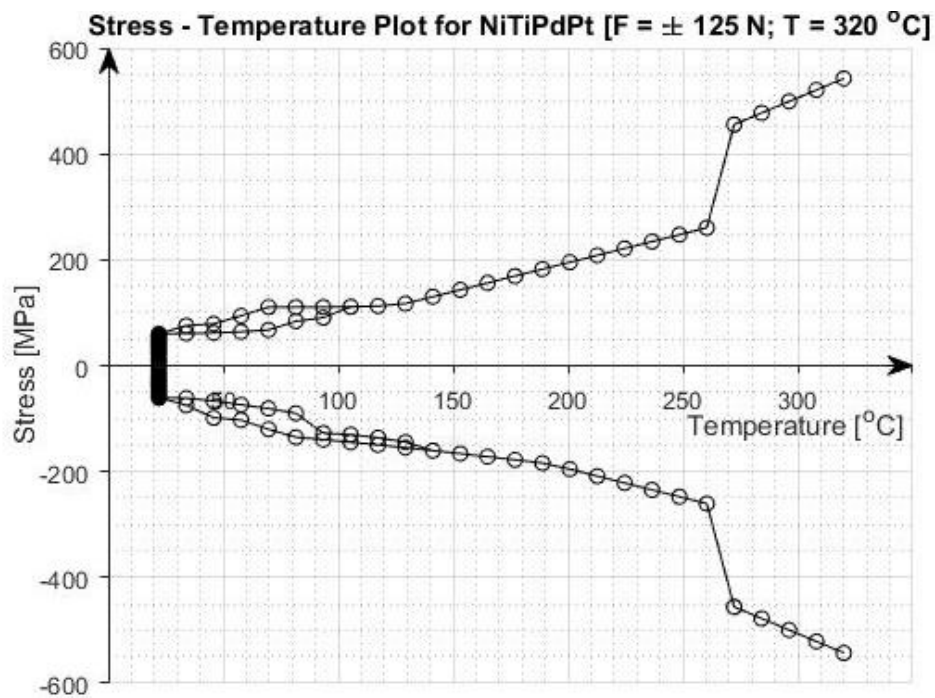
Figure 3.25 Stress vs. Force for both SMAs under  $F = \pm 100$  N.

### 3.2.3 NiTiPdPt under 125 N force

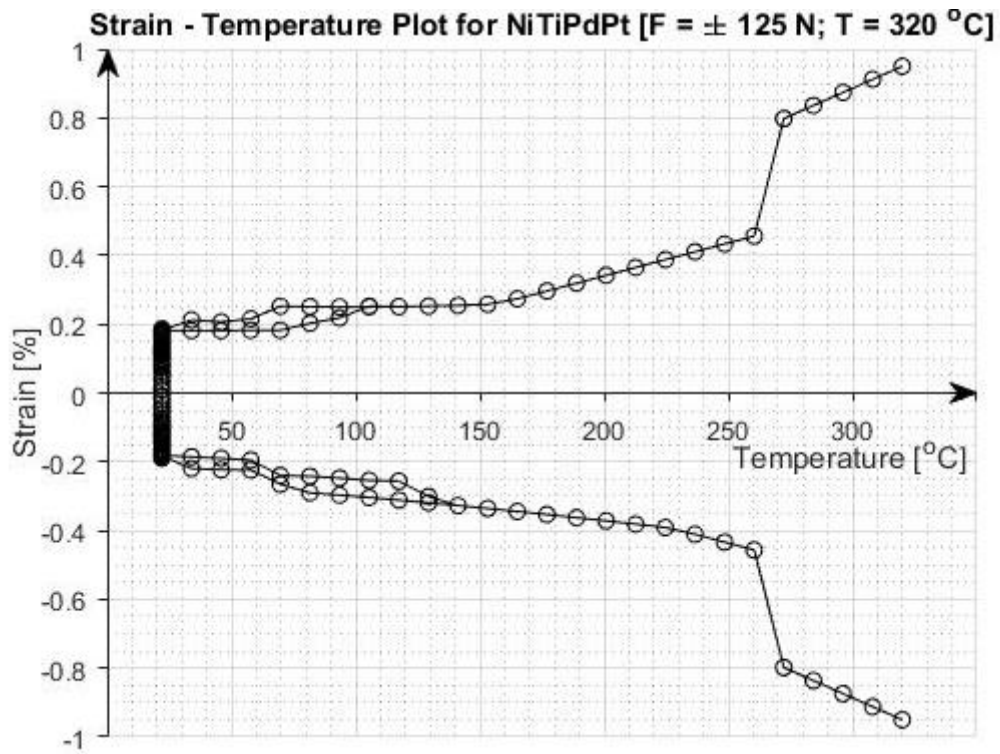
The following plots represent the results of Ni<sub>19.5</sub>Ti<sub>50.5</sub>Pd<sub>25</sub>Pt<sub>5</sub> at 320 °C under force of ± 125 N.



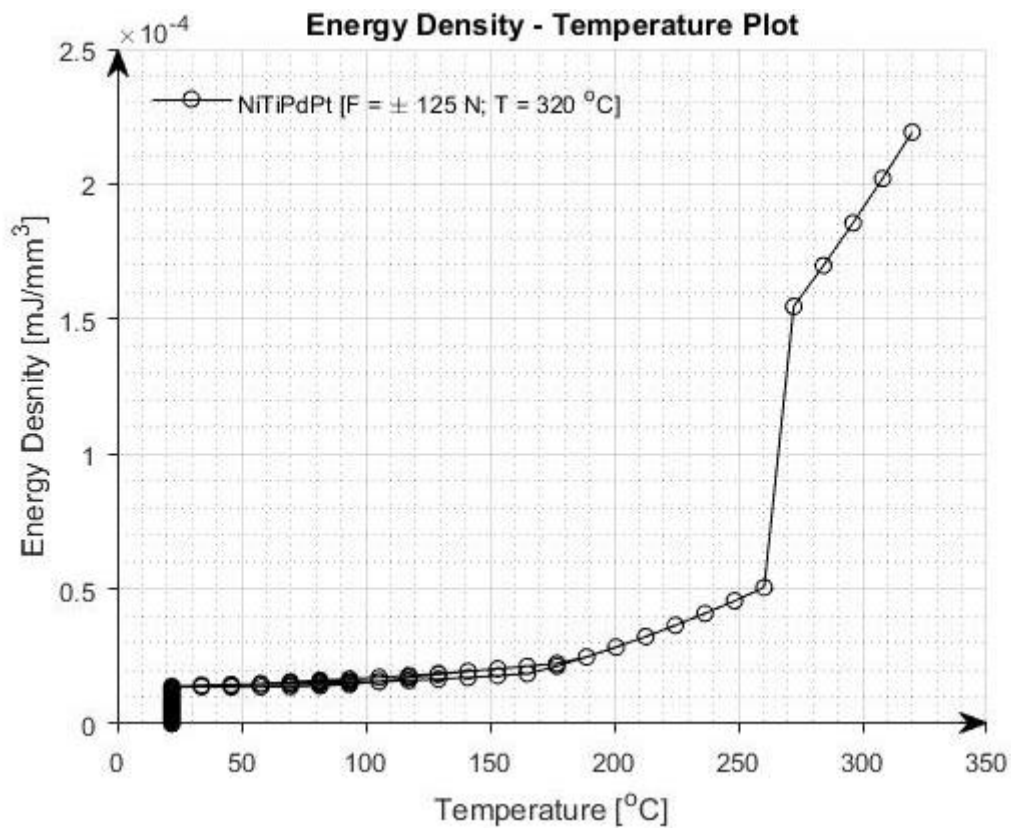
**Figure 3.26** Deflection vs. Temperature for NiTiPdPt at 320 °C under F = ± 125 N.



**Figure 3.27** Stress vs. Temperature for NiTiPdPt at 320 °C under F = ± 125 N.

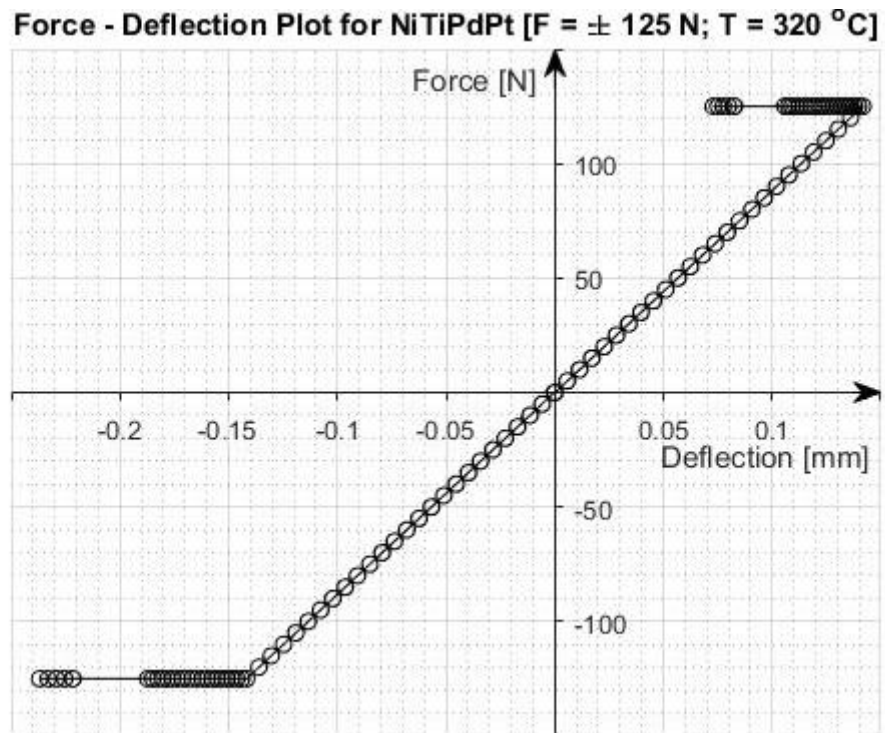


**Figure 3.28** Strain vs. Temperature for NiTiPdPt at 320 °C under  $F = \pm 125$  N.

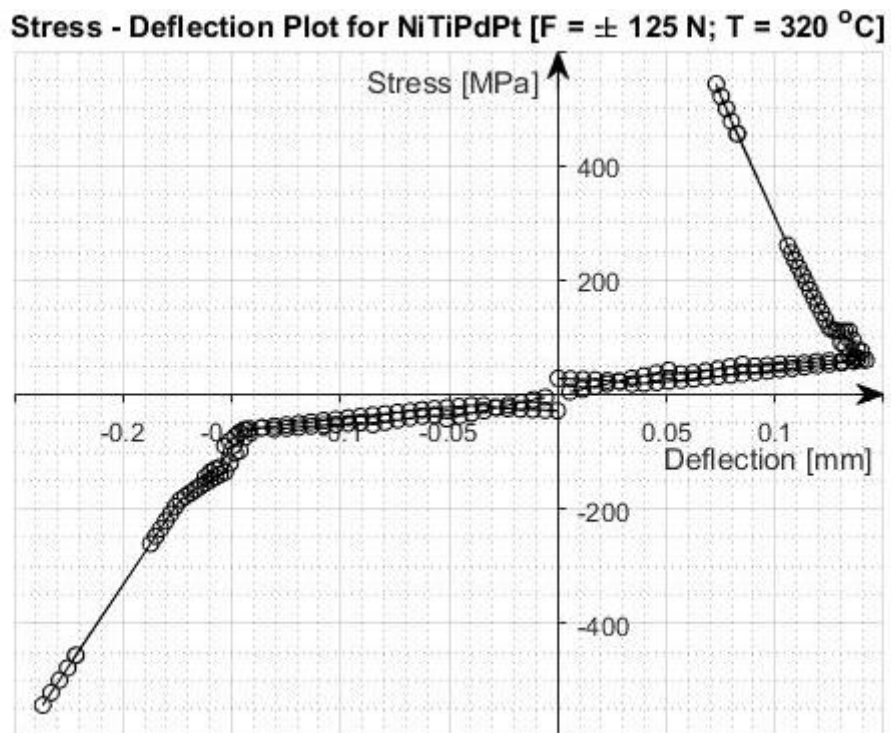


**Figure 3.29** Energy density vs. Temperature for NiTiPdPt at 320 °C under  $F = \pm 125$  N.

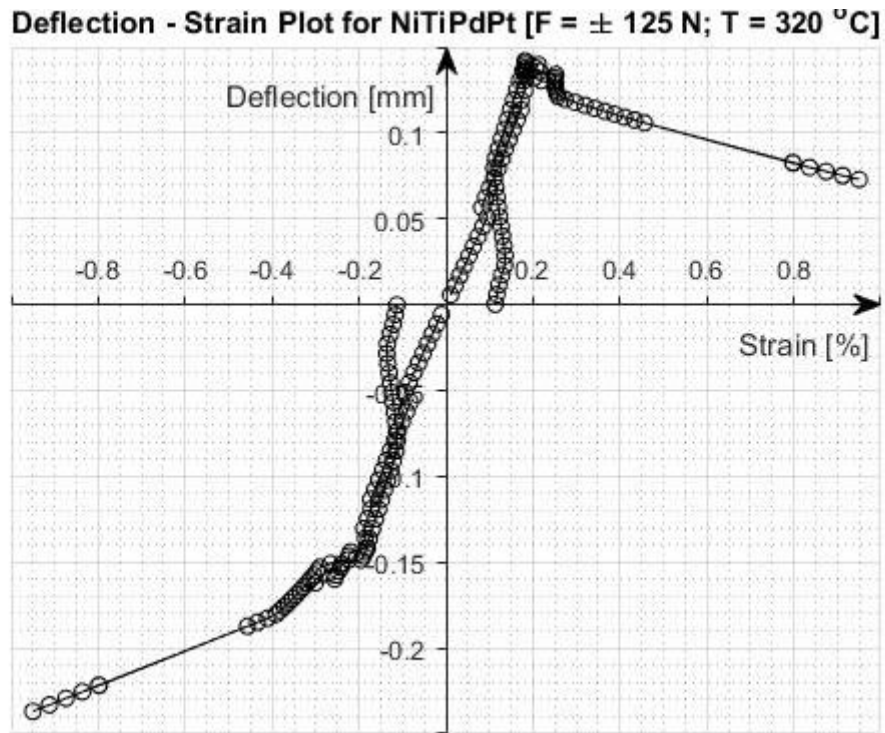




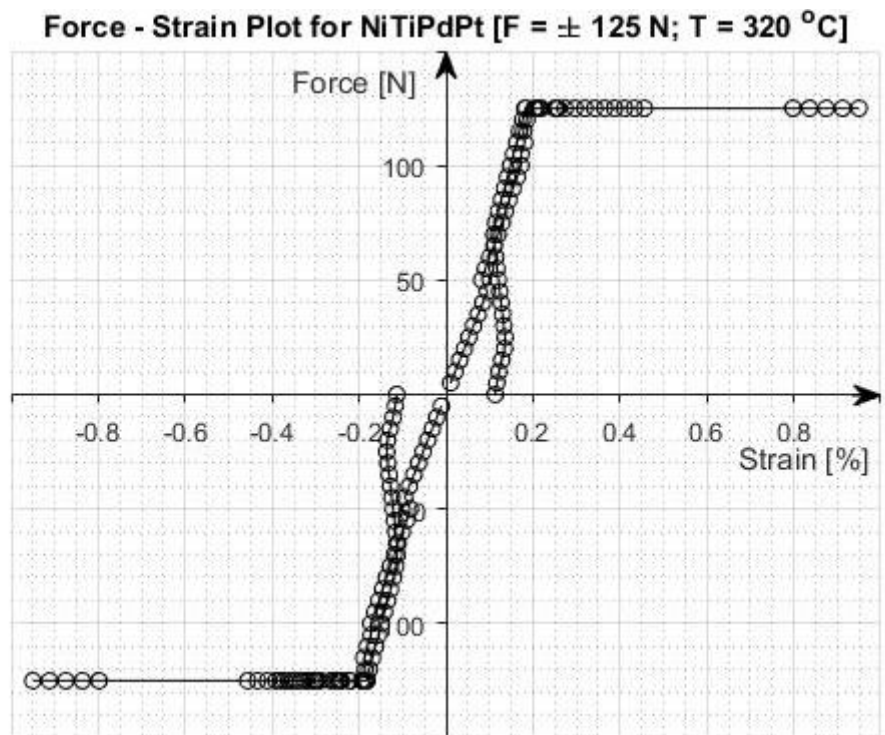
**Figure 3.30** Force vs. Deflection for NiTiPdPt at 320 °C under  $F = \pm 125 \text{ N}$ .



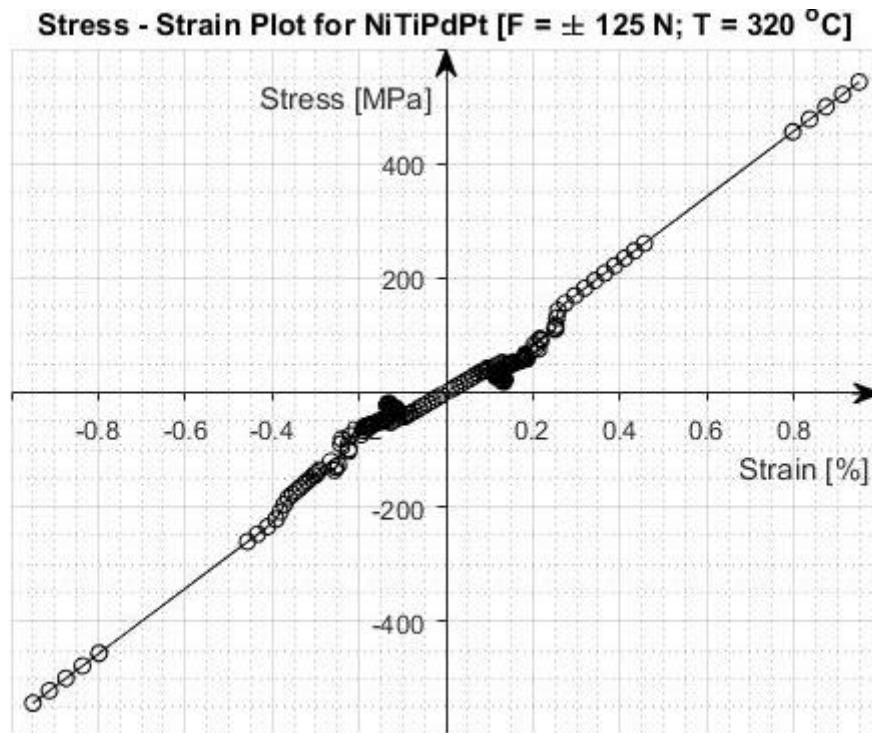
**Figure 3.31** Stress vs. Deflection for NiTiPdPt at 320 °C under  $F = \pm 125 \text{ N}$ .



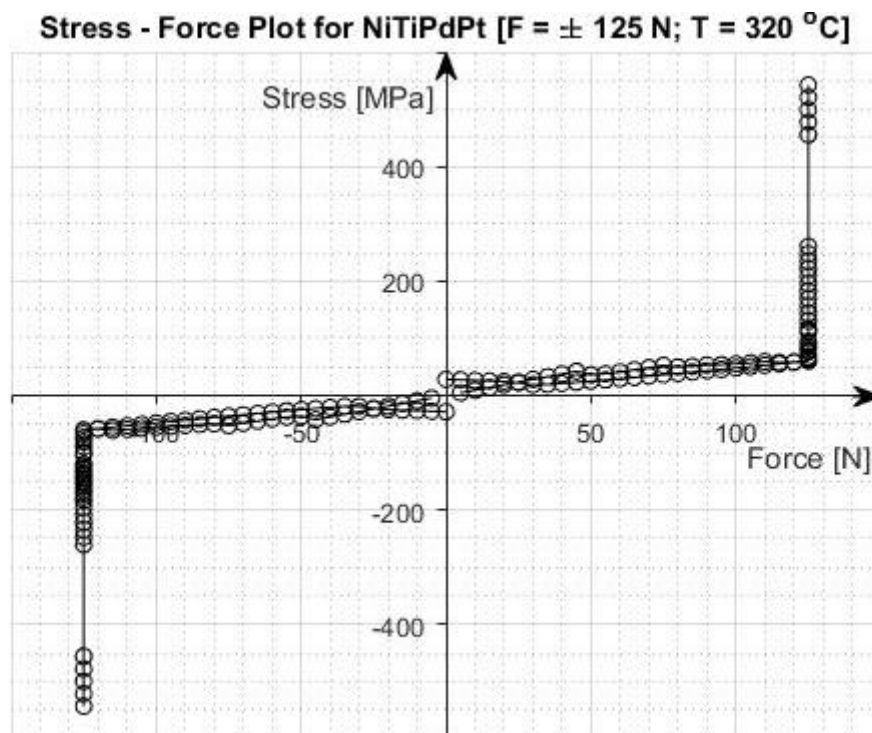
**Figure 3.32** Deflection vs. Strain for NiTiPdPt at 320 °C under  $F = \pm 125 \text{ N}$ .



**Figure 3.33** Force vs. Strain for NiTiPdPt at 320 °C under  $F = \pm 125 \text{ N}$ .



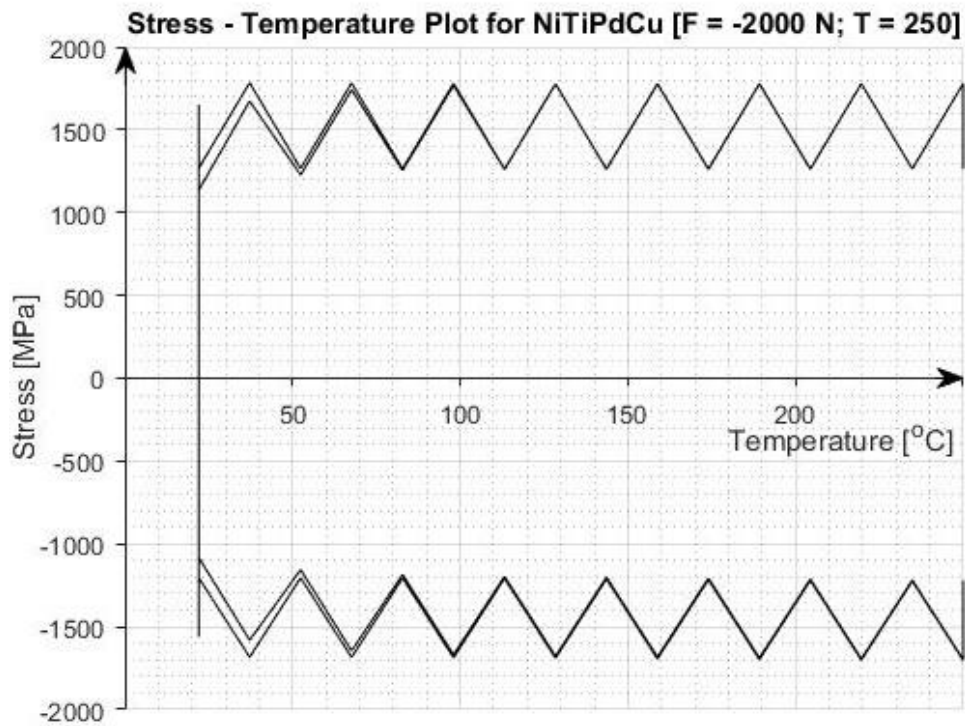
**Figure 3.34** Stress vs. Strain for NiTiPdPt at  $320 \text{ }^\circ\text{C}$  under  $F = \pm 125 \text{ N}$ .



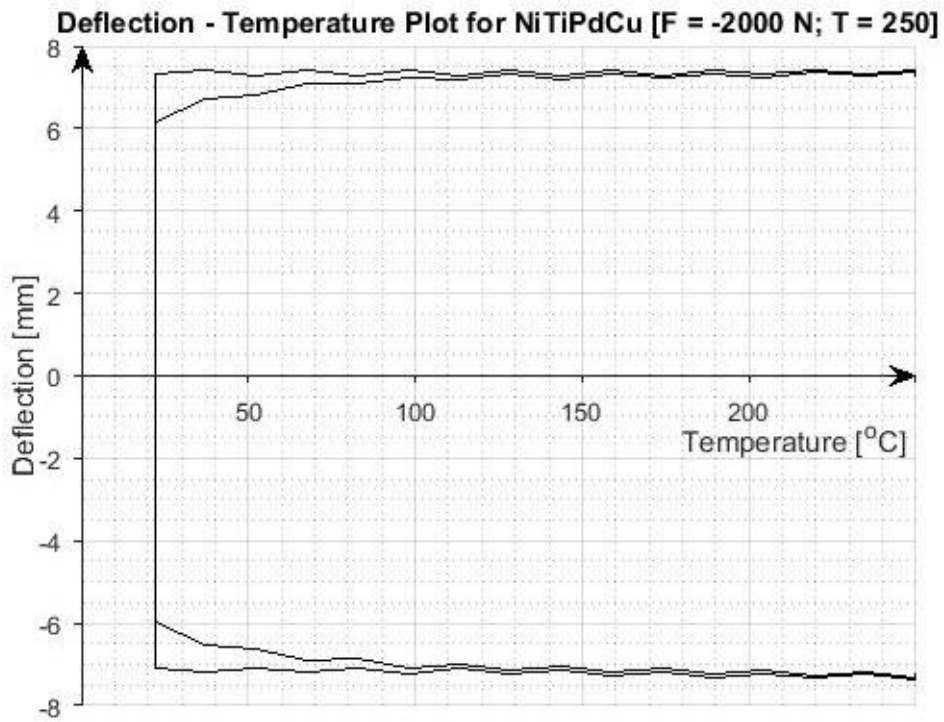
**Figure 3.35** Stress vs. Force for NiTiPdPt at  $320 \text{ }^\circ\text{C}$  under  $F = \pm 125 \text{ N}$ .

**3.2.4 NiTiPdCu under 2000 N force**

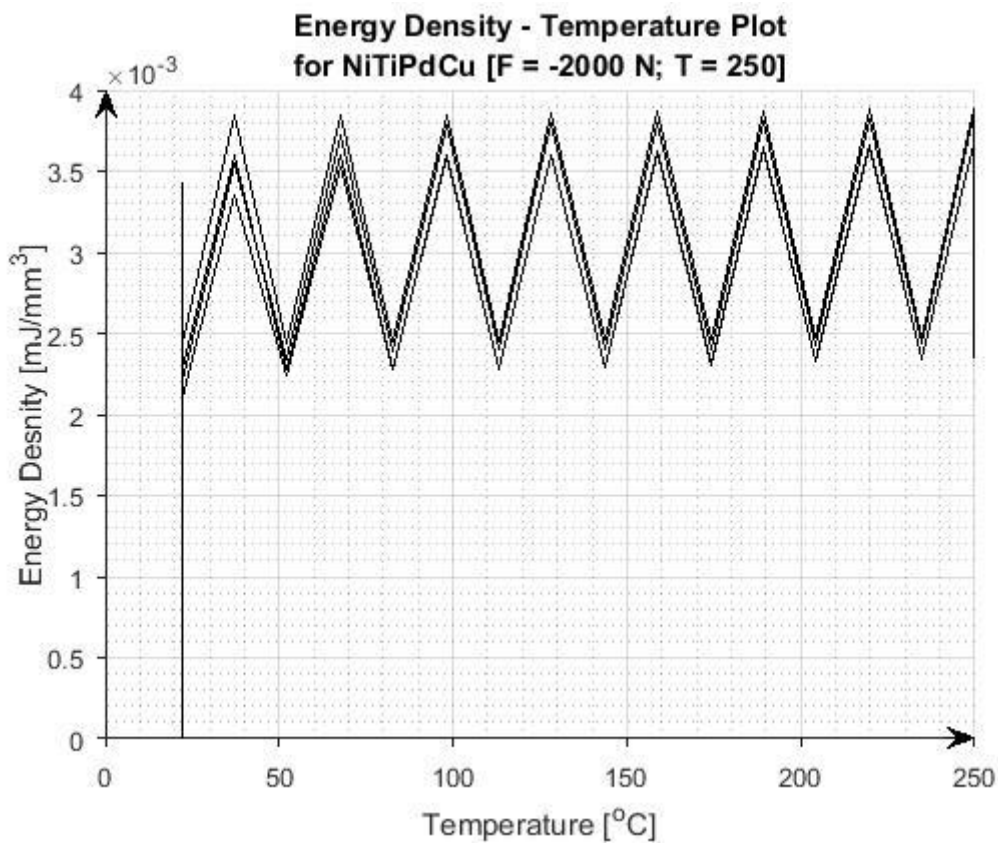
The following plots represent the results of Ni<sub>15</sub>Ti<sub>150</sub>Pd<sub>25</sub>Cu<sub>10</sub> at 250 °C under force of 2000N



**Figure 3.36** Stress vs. Temperature for NiTiPdCu at T = 250 °C under F = ± 2000 N.

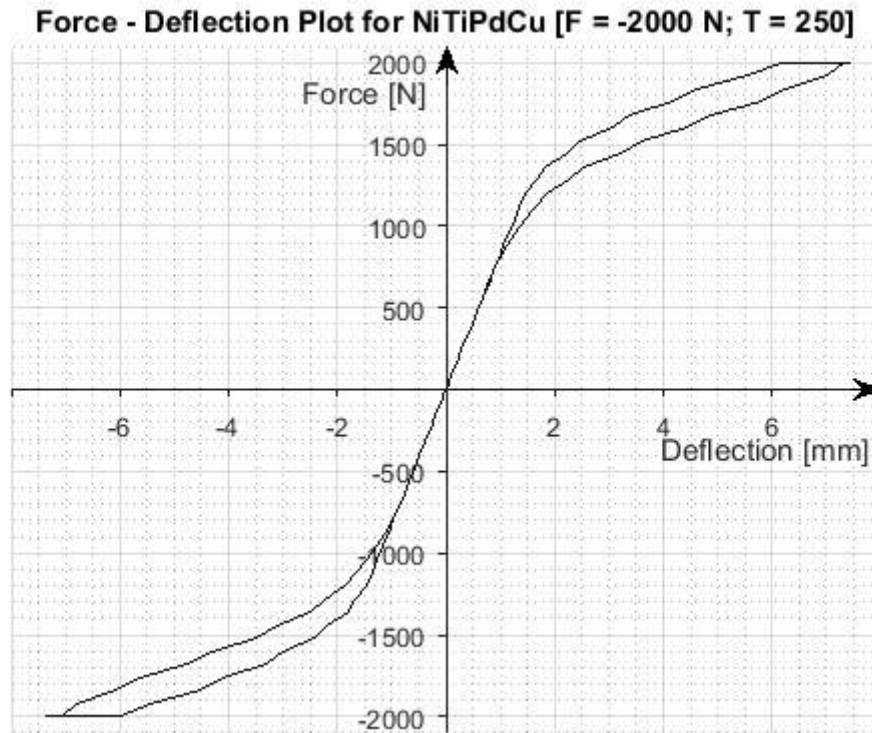


**Figure 3.37** Defection vs. Temperature for NiTiPdCu at T = 250 °C under F = ± 2000 N.

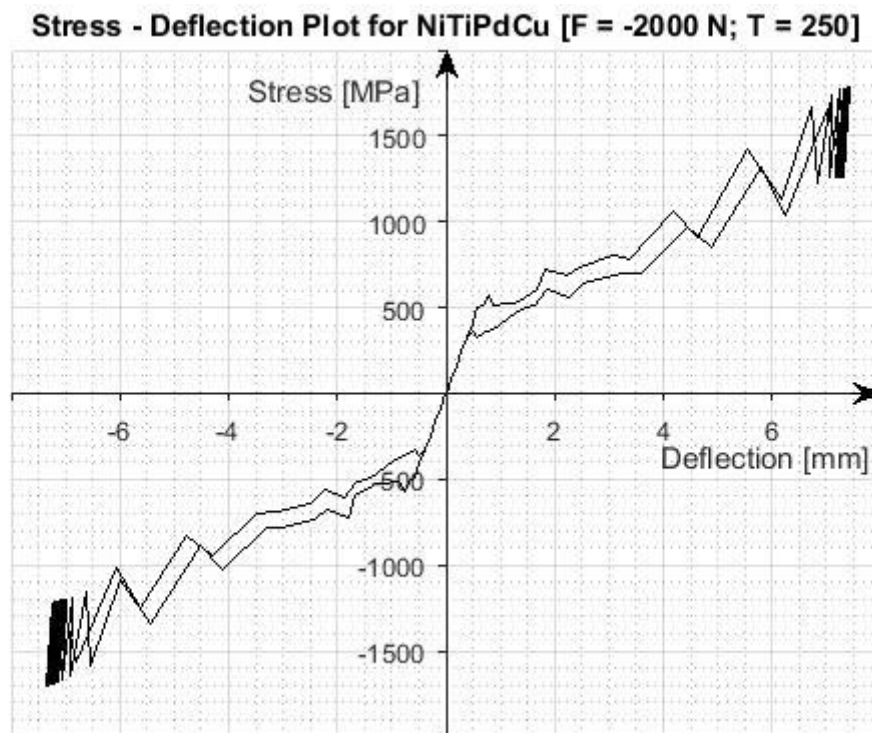


**Figure 3.38** Energy density vs. Temperature for NiTiPdCu at T = 250 °C under F = ± 2000

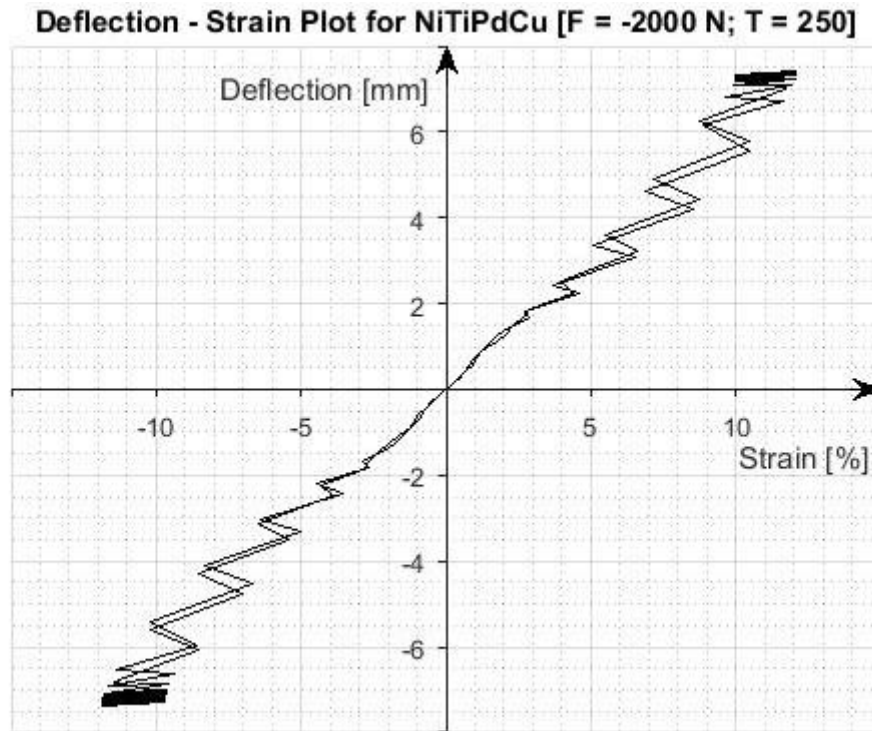
N.



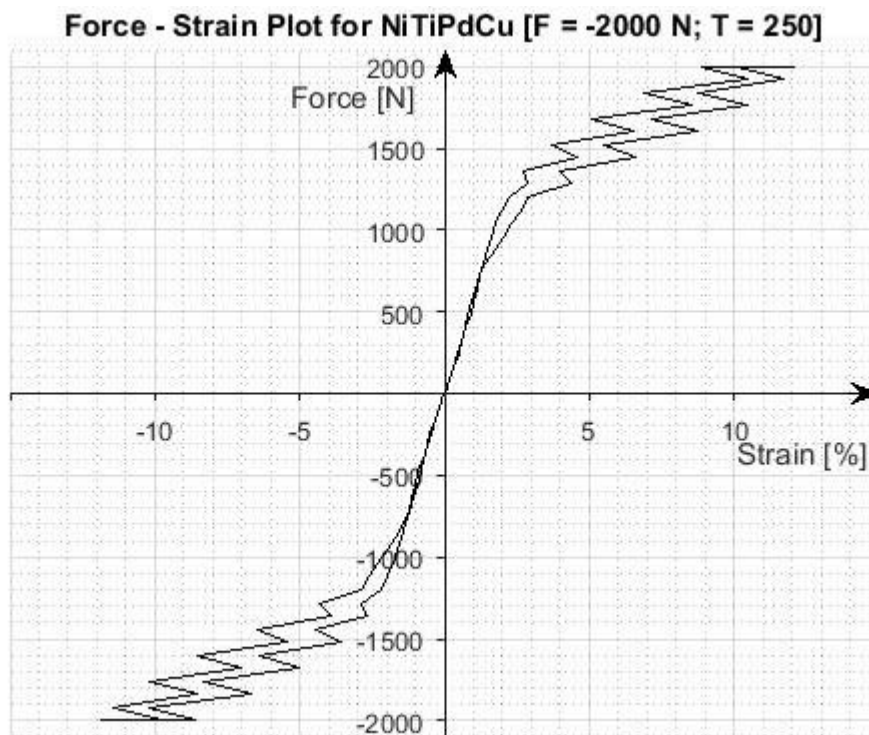
**Figure 3.39** Force vs. Deflection for NiTiPdCu at T = 250 °C under F = ± 2000 N.



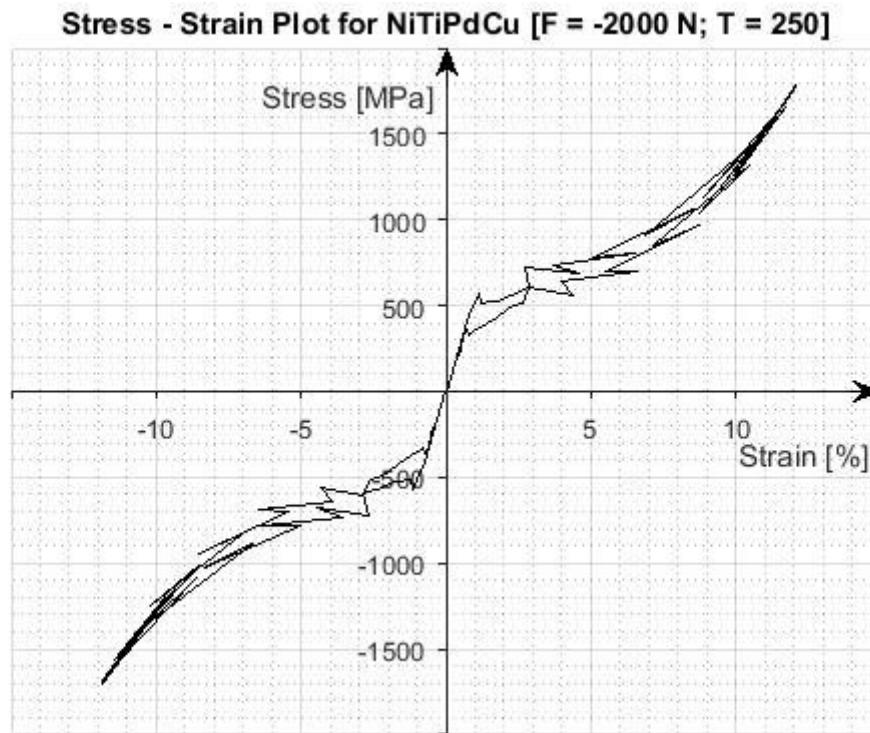
**Figure 3.40** Stress vs. Deflection for NiTiPdCu at T = 250 °C under F = ± 2000 N.



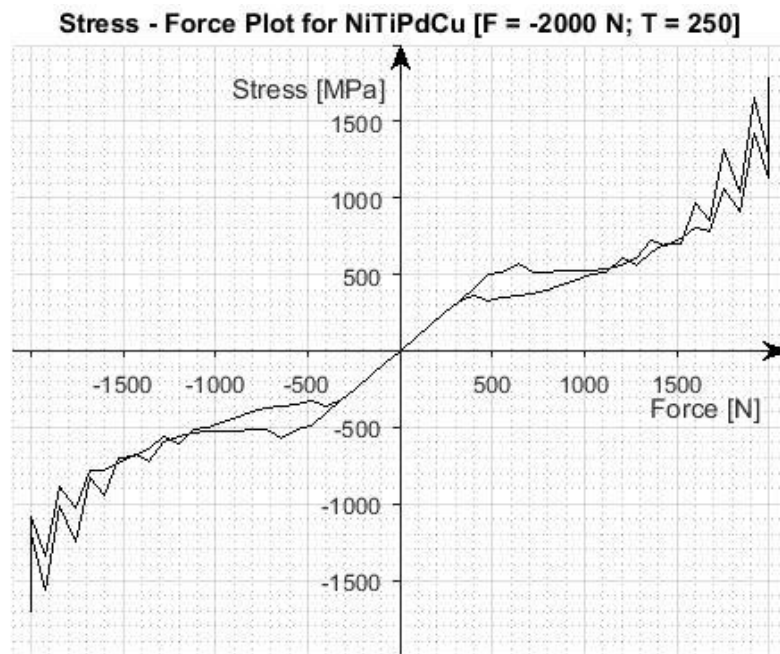
**Figure 3.41** Deflection vs. Strain for NiTiPdCu at  $T = 250\text{ }^{\circ}\text{C}$  under  $F = \pm 2000\text{ N}$ .



**Figure 3.42** Force vs. Strain for NiTiPdCu at  $T = 250\text{ }^{\circ}\text{C}$  under  $F = \pm 2000\text{ N}$ .



**Figure 3.43** Stress vs. Strain for NiTiPdCu at T = 250 °C under F = ± 2000 N.

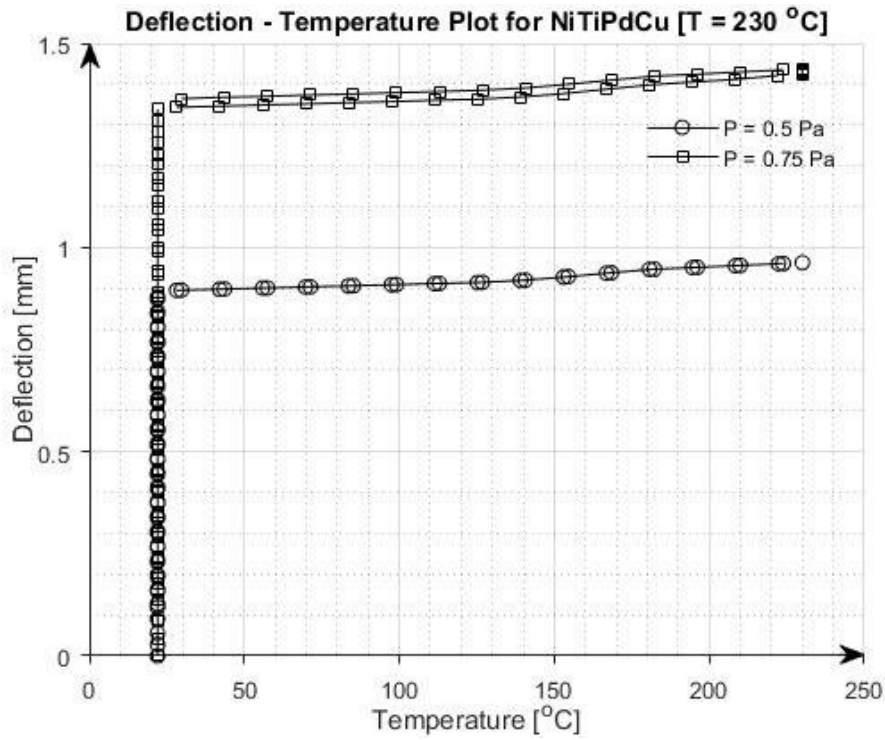


**Figure 3.44** Stress vs. Force for NiTiPdCu at T = 250 °C under F = ± 2000 N.

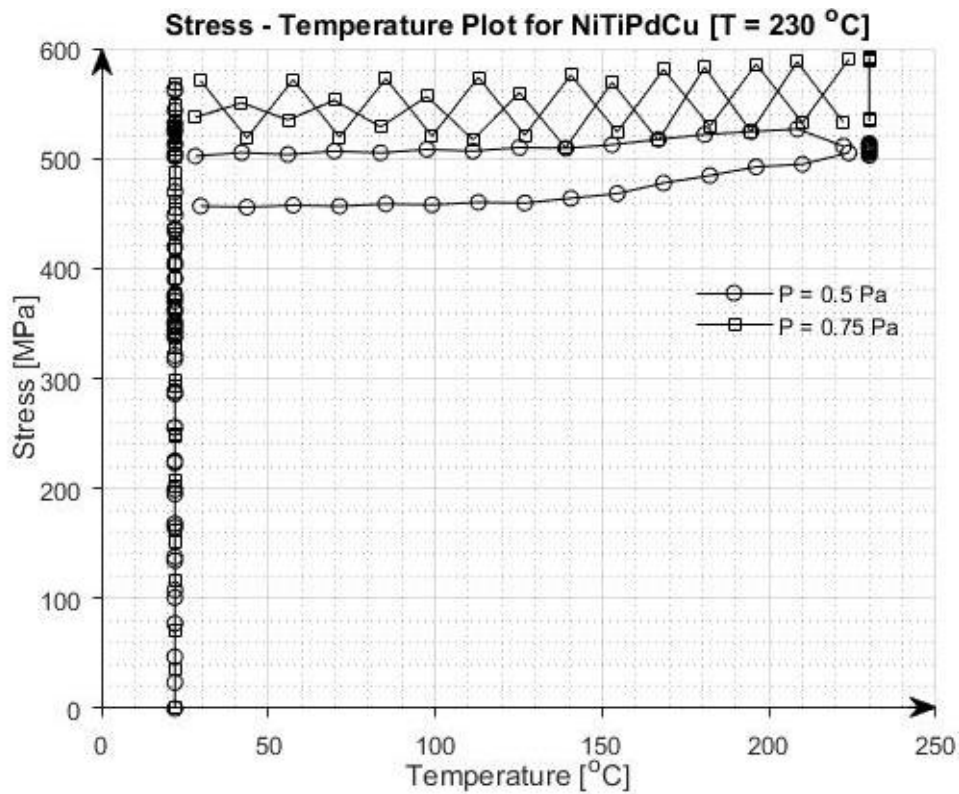
### 3.2.5 NiTiPdCu under 0.50 and 0.75 Pa pressure

The following plots represent the results of Ni<sub>15</sub>Ti<sub>50</sub>Pd<sub>25</sub>Cu<sub>10</sub> at 230 °C under pressure of 0.50 Pa and 0.75 Pa.

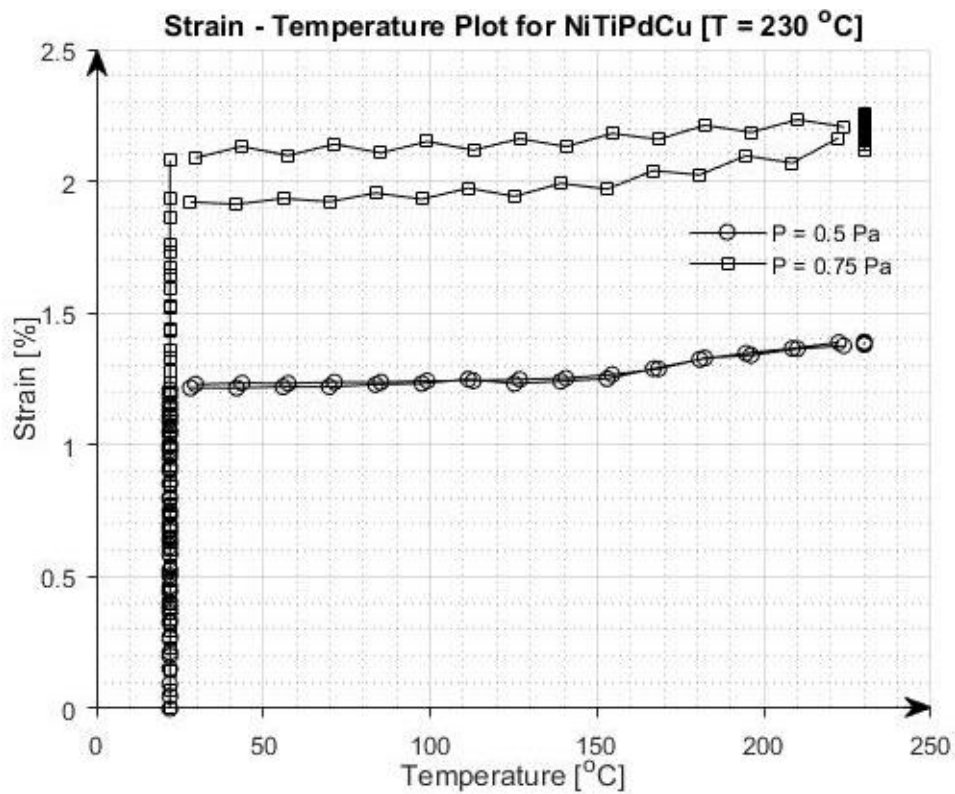




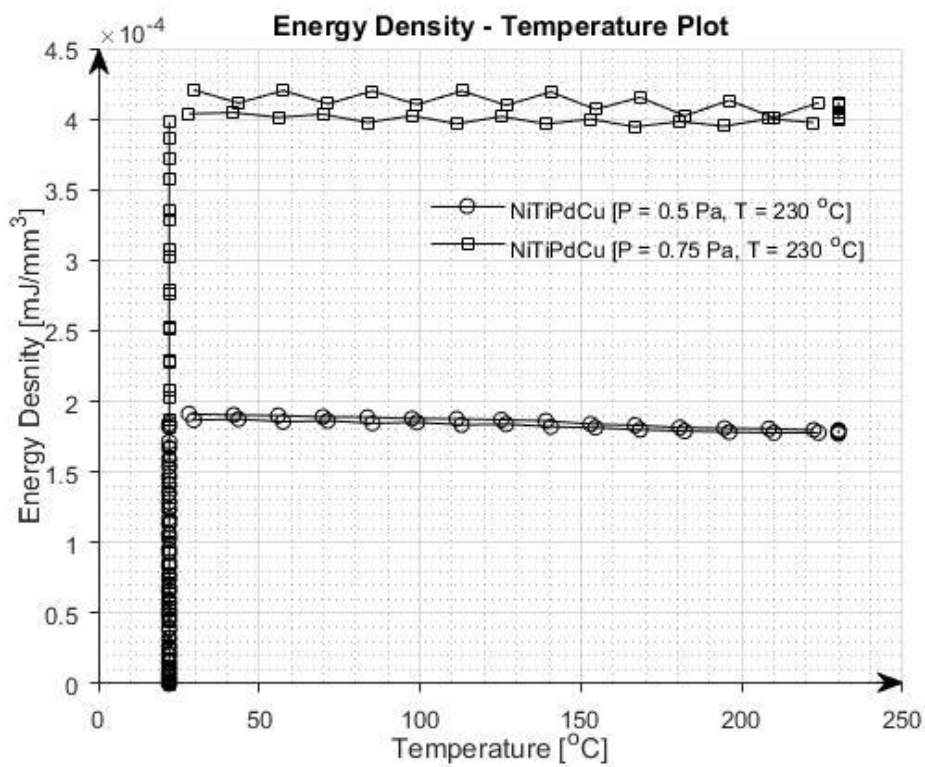
**Figure 3.45** Deflection vs. Temperature for NiTiPdCu at T = 230 °C.



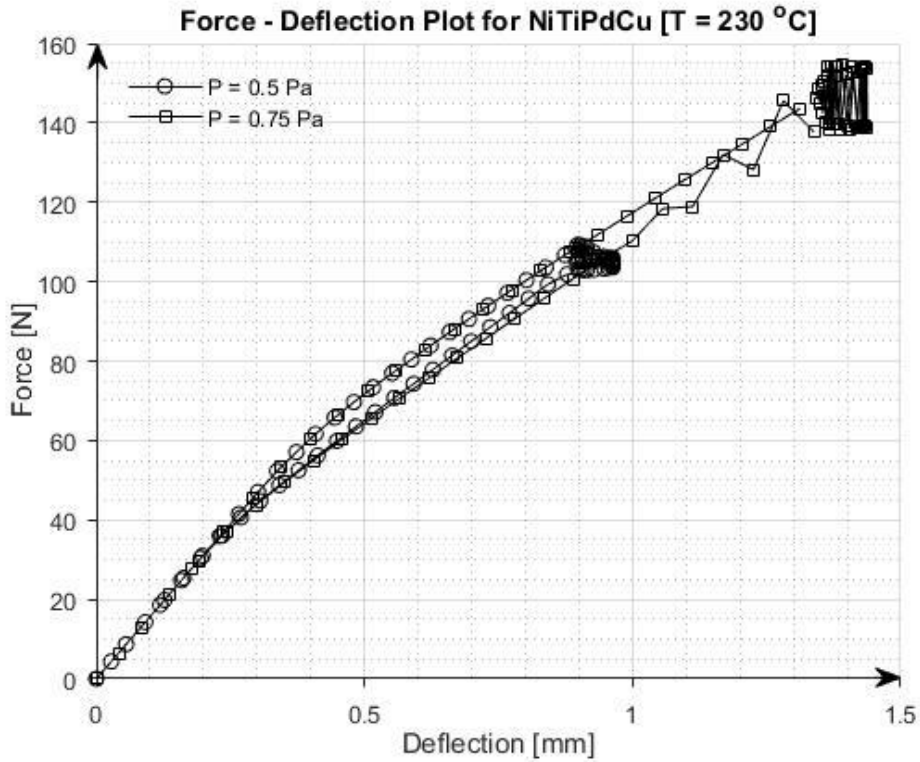
**Figure 3.46** Stress vs. Temperature for NiTiPdCu at T = 230 °C.



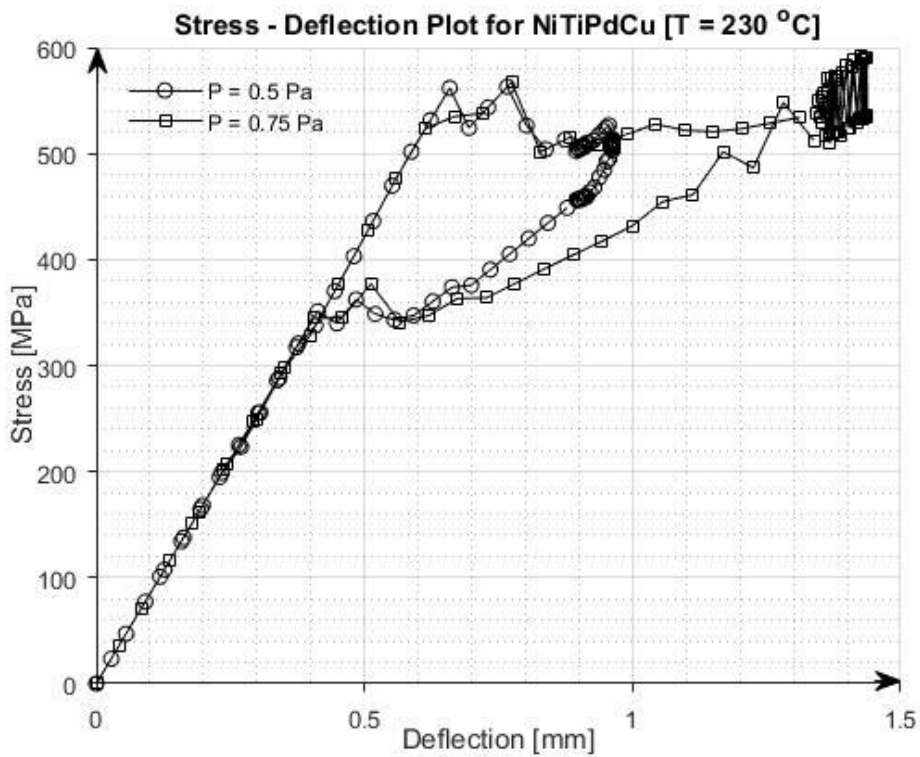
**Figure 3.47** Strain vs. Temperature for NiTiPdCu at T = 230 °C.



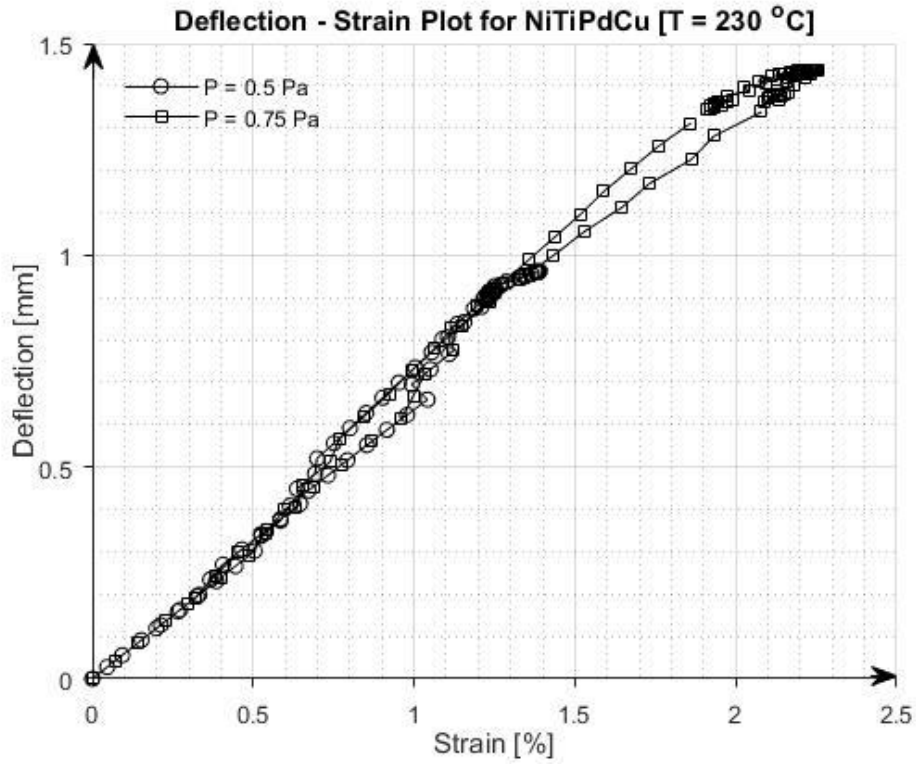
**Figure 3.48** Energy density vs. Temperature for NiTiPdCu at T = 230 °C.



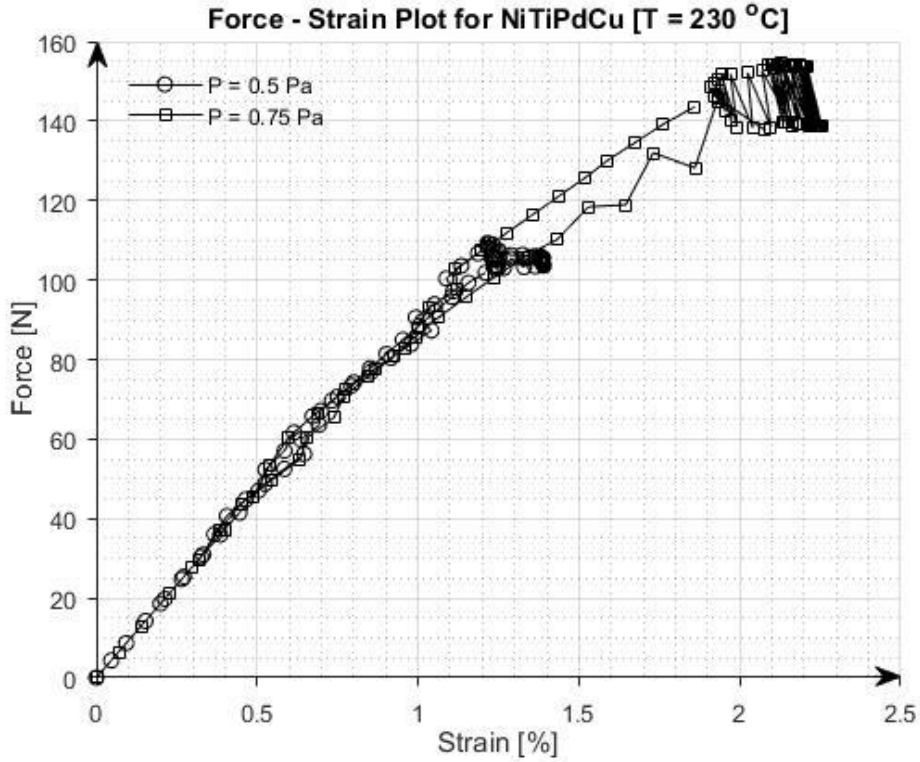
**Figure 3.49** Force vs. Deflectin for NiTiPdCu at T = 230 °C.



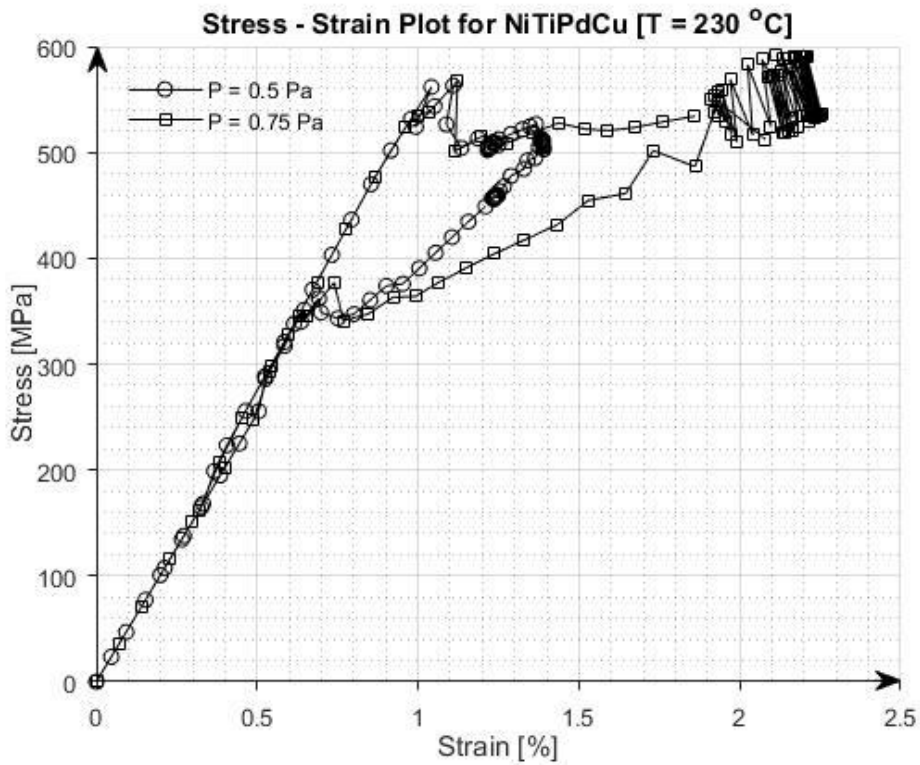
**Figure 3.50** Stress vs. Deflection for NiTiPdCu at T = 230 °C.



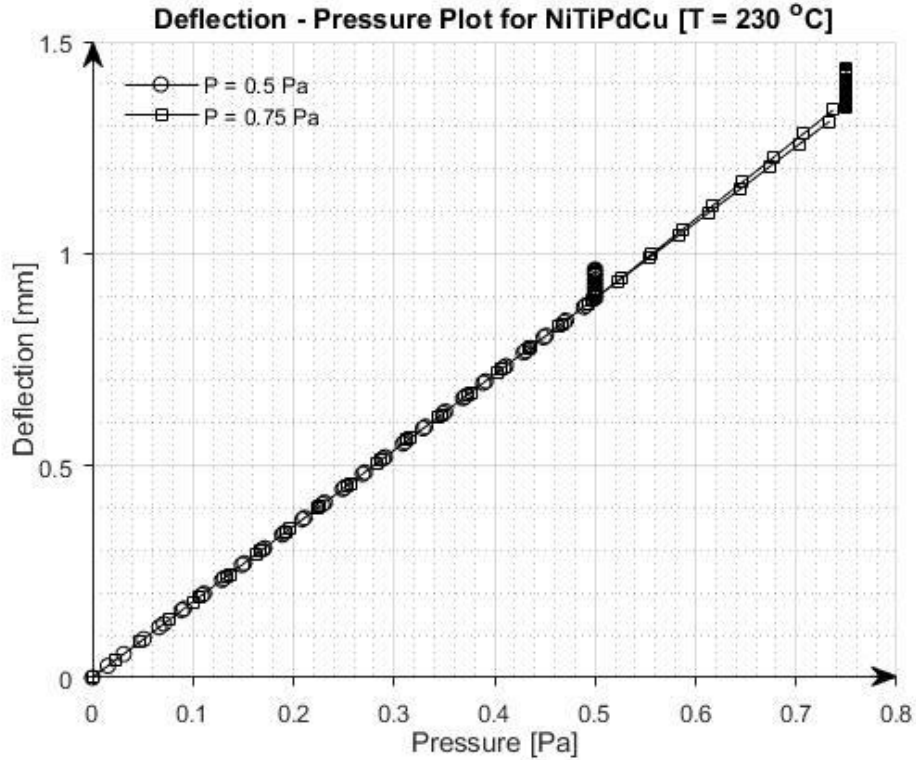
**Figure 3.51** Deflection vs. Strain for NiTiPdCu at  $T = 230 \text{ }^\circ\text{C}$ .



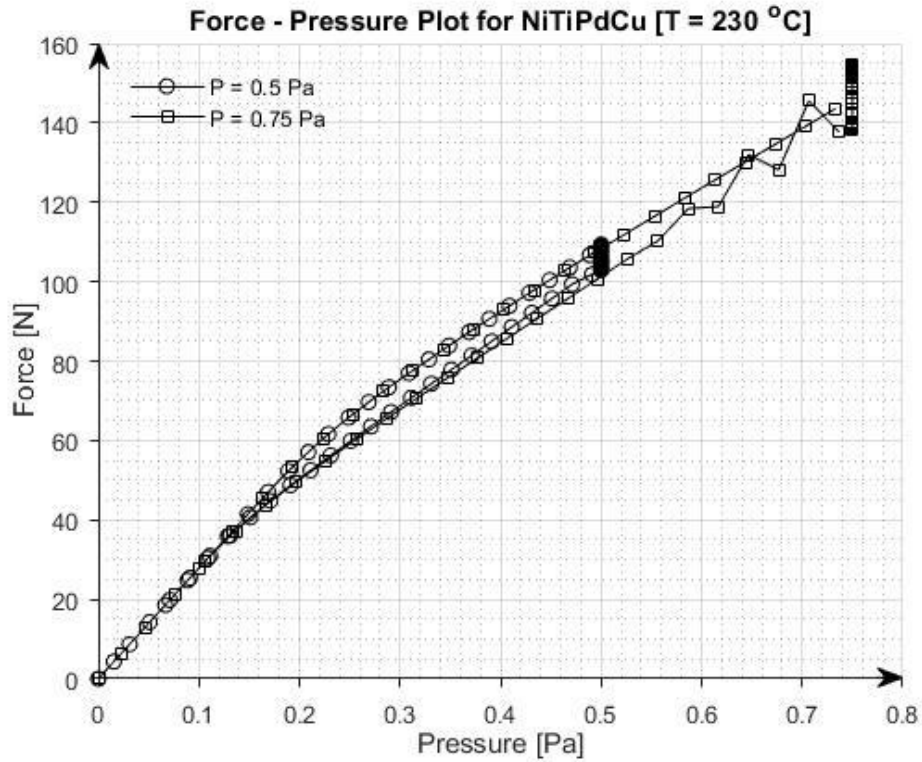
**Figure 3.52** Force vs. Strain for NiTiPdCu at  $T = 230 \text{ }^\circ\text{C}$ .



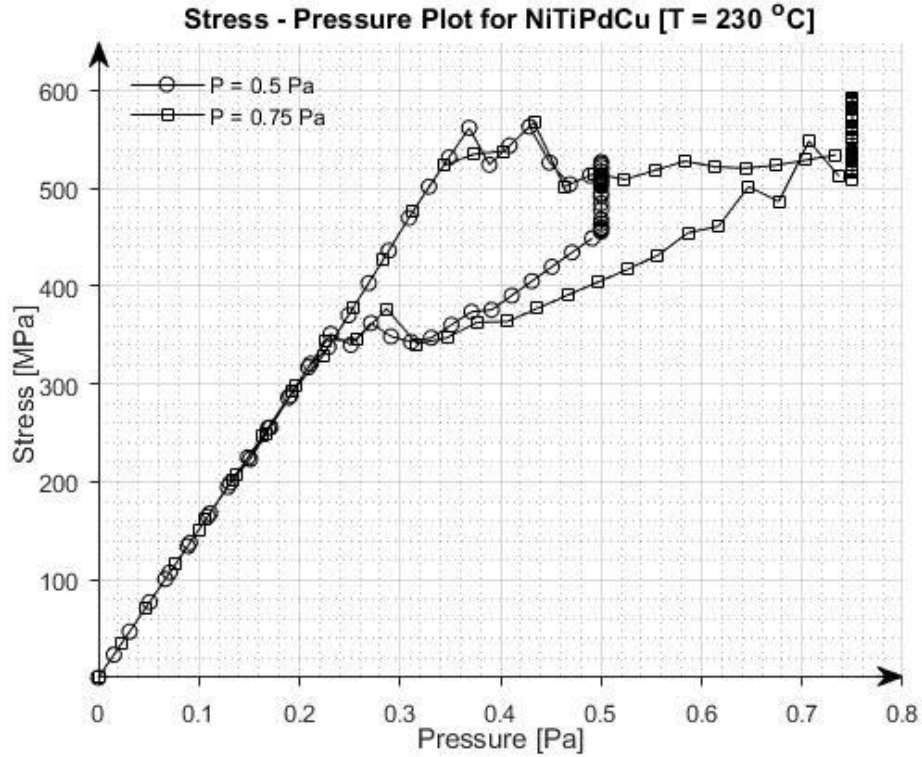
**Figure 3.53** Stress vs. Strain for NiTiPdCu at T = 230 °C.



**Figure 3.54** Deflection vs. Pressure for NiTiPdCu at T = 230 °C.



**Figure 3.55** Force vs. Pressure for NiTiPdCu at T = 230 °C.



**Figure 3.56** Stress vs. Pressure for NiTiPdCu at T = 230 °C.

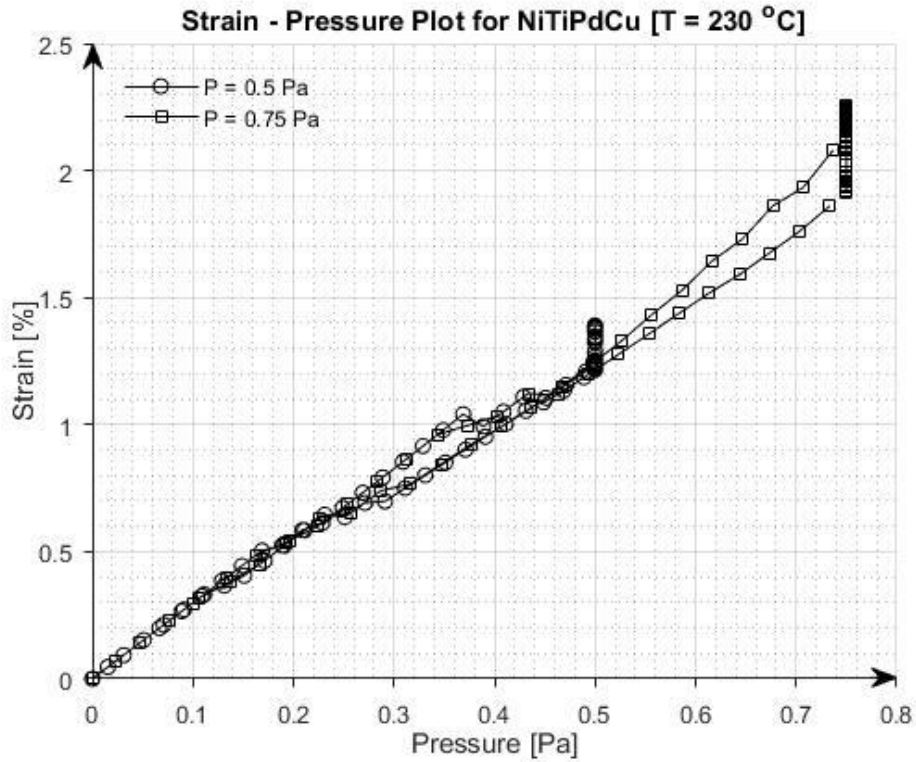


Figure 3.57 Strain vs. Pressure for NiTiPdCu at T = 230 °C.

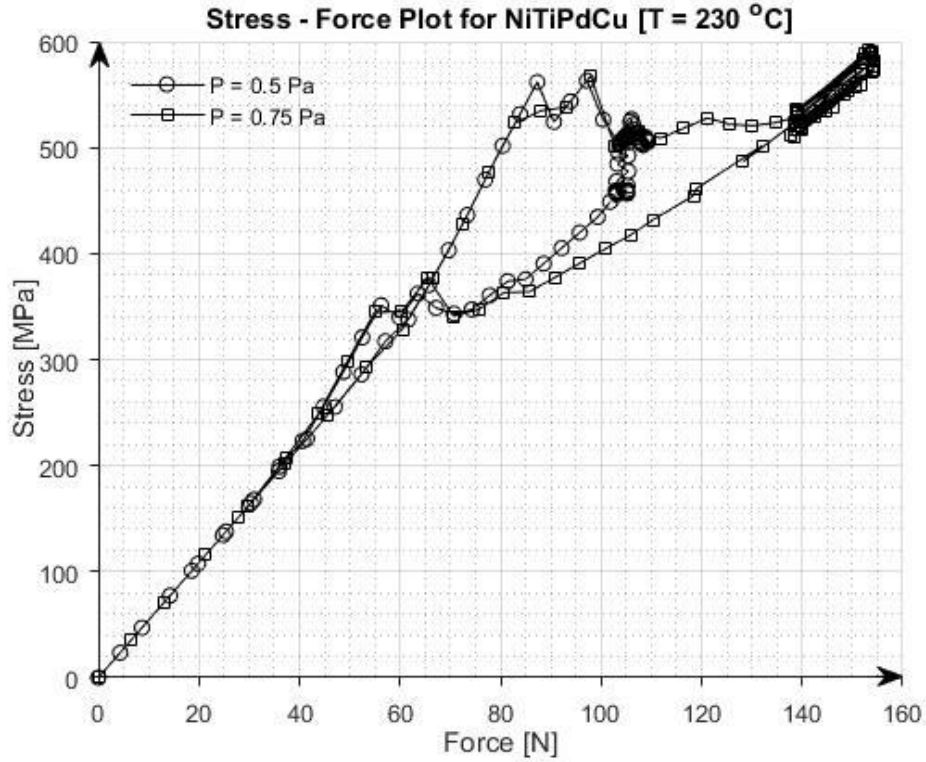
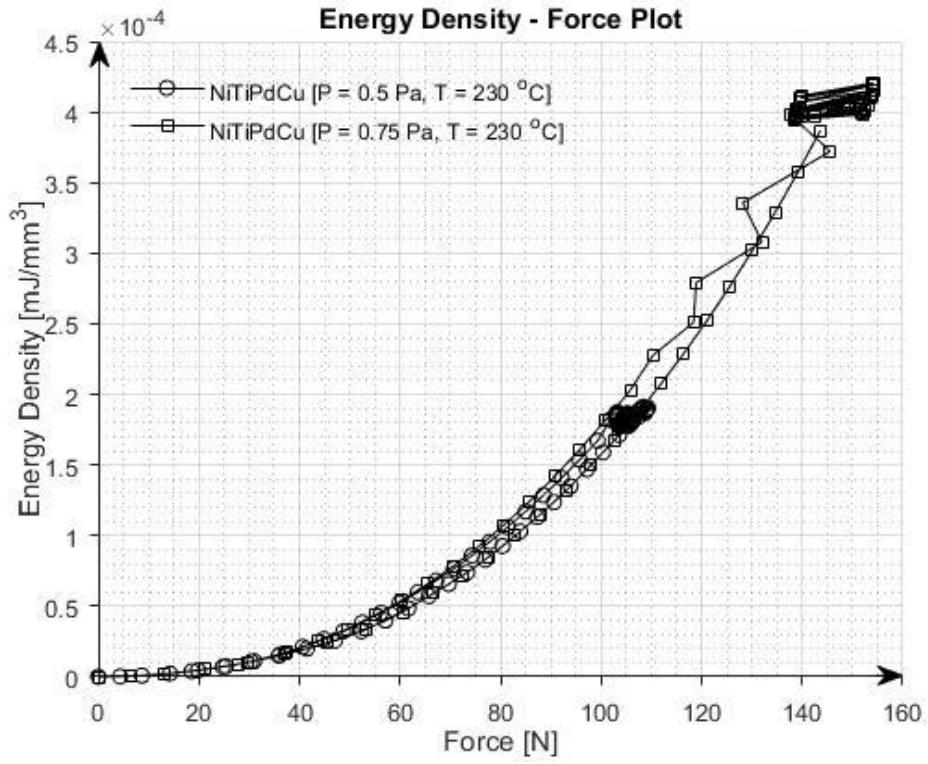


Figure 3.58 Stress vs. Force for NiTiPdCu at T = 230 °C.



**Figure 3.59** Energy density vs. Force for NiTiPdCu at T = 230 °C.



## CHAPTER 4

### RESULTS AND DISCUSSION

The spring is analysed with two smart materials  $\text{Ni}_{19.5}\text{Ti}_{50.5}\text{Pd}_{25}\text{Pt}_5$  and  $\text{Ni}_{15}\text{Ti}_{50}\text{Pd}_{25}\text{Cu}_{10}$  at various temperatures and under a range of axial forces. In simulations, the spring is analysed under  $\pm 40$  N,  $\pm 100$  N, 125 N, 2000 N force and under 0.50 and 0.75 Pa pressure. All these analyses are run at 250 °C and 320 °C. The following parameters are analysed in each simulation:

1. Deflection
2. Stress
3. Strain
4. Energy density

And these parameters are analysed against the following parameters:

1. Temperature
2. Deflection
3. Strain
4. Force

#### 4.1 Results

##### 4.1.1 SMAs under 40 N force

Figure 3.6 represents that both SMAs follows the same path of deflection against temperature plot and also return from the same path. Figure 3.7 represents that the path of both SMAs are different in stress against temperature plot although the both SMAs kept their path for return. The same behaviour is observed in the strain against temperature plot (Figure 3.8). Figure 3.9 represents that NiTiPdCu moved up and returned in the same pattern but NiTiPdPt did not follow the same path for return in energy density against temperature plot. Figure 3.10

represents that both SMAs are very close in following the curve for force against deflection plot though when the spring is compressed, the return path is different.

Figure 3.11 represents the same behaviour with a change that appears in NiTiPdPt which shows too high stress as compare to NiTiPdCu under  $\pm 40$  N force in stress against deflection plot. Figure 3.12 represents the deflection of NiTiPdCu stopped at 0.6% while that of NiTiPdPt went up to 0.95% though both SMAs return following the same path in deflection against strain plot. Figure 3.14 represents the straight line plot without any sign of transformation of both SMAs under  $\pm 40$  N force in stress against strain plot. Figure 3.15 represents that the effect of force on spring is minimal i.e. only 40 MPa stress was built up under the force of 39 N in stress against force plot.

#### **4.1.2 SMAs under 100 N force**

Figure 3.16 represents the different paths followed by each SMA whereas on return, NiTiPdCu took the same path while NiTiPdPt deviated a bit from its original path in deflection against temperature plot. The same behaviour is observed for the both SMAs in stress against temperature plot (Figure 3.17). Figure 3.18 represents that the deviation is captured at transformation phase only, rest of the curve is same for both ways in strain against temperature plot for both SMAs. The same behaviour is again observed in energy density against temperature plot (Figure 3.19).

Figure 3.20 represents the path followed by both SMAs are same in force deflection plot. Figure 3.21 represents the difference of paths of both SMAs, the stresses reached up to 550 MPa in case of NiTiPdPt while it remained 325 MPa in case of NiTiPdCu in stress deflection plot. Figure 3.22 represents the same abrupt behaviour of path for NiTiPdPt though NiTiPdCu followed the same path in both ways in deflection against deflection plot.

Figure 3.23 represents the residual strain on the removal of force from NiTiPdPt but no such occurrence observed in NiTiPdCu in force against strain plot. Figure 3.24 represents the straight line behaviour without any phase transformation in both SMAs in stress against strain plot. Figure 3.25 represents more inclination of graph of NiTiPdCu than that of NiTiPdPt but the former has less stress than later in stress against force plot.

#### 4.1.3 NiTiPdPt under 125 N force

Figure 3.26 represents the same path followed for both ways for NiTiPdPt in deflection against temperature plot. Figure 3.27 represents the same path followed at the transformation temperature range but the different of path is observed at 25 °C to 125 °C temperature in stress against temperature plot. The same behaviour is observed in strain against temperature plot (Figure 3.28). Figure 3.29 represents the same path followed for both ways in transformation phase but it shows deviation at 50 °C to 160 °C in energy density against temperature plot.

Figure 3.30 represents the same path followed in both ways while applying force or removal of force in force against deflection plot. Figure 3.31 represents a bit abruptness in the behaviour of both ways but that small deviation can be neglected on the basis of alloying composition in stress against deflection plot. Figure 3.32 represents the effect of residual strain up to 0.1% for NiTiPdPt in deflection against strain plot. The same behaviour is observed in force against strain plot (Figure 3.33). Figure 3.34 represents the same straight line behaviour which followed the same path in both ways in stress against strain plot. Figure 3.35 represents a small amount of residual stresses on the removal of force i.e. about 15 MPa though it followed the same path until on removal of force up to 50 N in stress against force plot.

#### 4.1.4 NiTiPdCu under 2000 N force

Figure 3.36 represents the same path followed behaviour on heating and cooling with a small deflection which reconciles on further cooling even under 2000 N compressive force in stress against temperature plot. The same behaviour is observed in deflection against temperature plot (Figure 3.37). Figure 3.38 represents a lot of energy density with much deviation in energy density against temperature plot. Figure 3.39 represents the typical behaviour of any SMA on application and removal of force and it shows hysteresis loop for NiTiPdCu in force against deflection plot.

Figure 3.40 represents the stress up to 1.8 GPa which returned on the different path but matches before removal of stress in stress against deflection plot. Figure 3.41 represents a small deviation of paths on extension and compression but it was removed in the absence of forces in deflection against strain plot. The same graph was observed in force against strain plot (Figure 3.42). Figure 3.43 represents the stress reaches up to 1.8 GPa while strain reaches up to 12% which was not achieved before with NiTiPdPt in stress against strain plot. The same behaviour is seen in stress against force plot (Figure 3.44).

#### 4.1.5 NiTiPdCu under 0.50 Pa and 0.75 Pa pressure

Figure 3.45 represents the deflection up to 0.8 mm and 1.35 mm for NiTiPdCu under pressure of 0.50 Pa and 0.75 Pa respectively in deflection against temperature plot. Figure 3.46 represents the stress of NiTiPdCu reaches up to 500 MPa and 580 MPa under pressure of 0.50 Pa and 0.75 Pa respectively in stress against temperature plot. Figure 3.47 represents the strain up to 1.4% and 2.2% under pressure of 0.50 Pa and 0.75 Pa respectively in strain temperature plot. Figure 3.48 represent notable difference in energy density i.e.  $2 \times 10^{-4}$  mJ/mm<sup>3</sup> and  $4.3 \times 10^{-4}$  mJ/mm<sup>3</sup> under pressure of 0.50 Pa and 0.75 Pa respectively in energy density against temperature plot.

Figure 3.49 represent the force reached up to 110 N and 155 N under pressure of 0.50 Pa and 0.75 Pa respectively in force deflection plot. Figure 3.50 represents much deviated graphs of stress which reconciles on the removal of pressure in stress and deflection plot. Figure 3.51 represents the behaviour is too close under pressure of 0.50 Pa and 0.75 Pa in deflection against strain plot. The same behaviour is observed in the force against strain plot (Figure 3.52). Figure 3.53 represents the straight line behaviour of NiTiPdCu under pressure of 0.50 Pa and 0.75 Pa in deflection against pressure plot. Figure 3.55 represents a small deviation in paths under pressure of 0.50 Pa and 0.75 Pa in force against pressure plot.

Figure 3.56 represents a huge deviation of paths and hysteresis loop under pressure of 0.50 Pa and 0.75 Pa in stress against pressure plot. Figure 3.57 represents very close paths of strains under pressure of 0.50 Pa and 0.75 Pa in strain against pressure plot. Figure 3.58 represent the same abrupt behaviour of stress which later on reconciles on the removal of pressure under 0.50 Pa and 0.75 Pa pressure in stress against deflection plot. Figure 3.59 represents the parabolic relationship of the curve under pressure of 0.50 Pa and 0.75 Pa in energy density against force plot.

## 4.2 Discussion

The behaviour of stress, strain, deflection and energy density against temperature, force and pressure are presented in the result section which gives very clear comparison of the both SMAs. In each graph plotted above in chapter 3 and as the results presented in this chapter, both SMAs show some advantage on one another but these positive points needs to be analyzed and discussed critically before reaching to a conclusion. The ability to withstand with force

without causing residual stresses in evident in NiTiPdCu while the effect of deflection is less appeared in NiTiPdPt. At some points, the curves of the plots deviated a lot from their original path and this behaviour is portrayed by both SMAs. Another important factor which decides the usability of a material as a spring material is energy density that how much energy can it absorb before failure. The more energy a material can absorb, the better it is considered for spring or actuator application. And also, how much energy is required to cause a change in unit length which can be observed in the form of strain.

## CHAPTER 5

### CONCLUSION AND FUTURE WORK

The results and discussion favours that the newly composed  $\text{Ni}_{15}\text{Ti}_{50}\text{Pd}_{25}\text{Cu}_{10}$  alloy is better for spring material than  $\text{Ni}_{19.5}\text{Ti}_{50.5}\text{Pd}_{25}\text{Pt}_5$  alloy since it has better thermomechanical properties than old one.

The curves and behaviour of numerically simulated NiTiPdCu is same as numerically and experimentally tested NiTiPdPt. This is evident for the soundness of the FEA to predict the behaviour of newly composed SMA rather than directing towards costly and time-consuming long way of experimentations. Before implementation in any FEA package, one must know the constitutive models available in that open-source or commercial FEA package. Unawareness may lead to drastically wrong results.

Since these simulations are run and analysed only under dynamic axial loading but torsional loading is also very crucial in determining the strength, life and fatigue of the spring. It was not the case here as only axial loading was present in the real case scenario hence to reduce computational power, only spring is tested as various temperatures and axial loadings. One future work could be added to simulate the same for torsional loading and later on superposition of both axial and torsional effects can be obtained for those scenarios where both loadings are present.

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